

# Chlorine Oxide Radicals $\text{ClO}_x$ ( $x = 1-4$ ) Studied by Matrix Isolation Spectroscopy

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**Abstract:** Low pressure flash thermolysis of different precursor molecules containing  $-\text{ClO}$ ,  $-\text{ClO}_3$  or  $-\text{OClO}_3$  yield, when highly diluted in Ne or  $\text{O}_2$  and subsequent quenching of the products in a matrix at 5 or 15 K,  $\text{ClO}_x$  ( $x = 1, 3, 4$ ) radicals, respectively. If Ne or  $\text{O}_2$  gas is directed over solid  $\text{ClO}_2$  at  $-120^\circ\text{C}$  and the resulting gas mixtures are immediately deposited as a matrix, a high fraction of  $(\text{OClO})_2$  is trapped. This enables recording of IR and UV spectra of weakly bonded  $(\text{OClO})_2$  dimers and detailed studying of their photochemistry. For Ne or  $\text{O}_2$  matrix isolated ClO radicals the vibrational wavenumbers and electronic transitions are only slightly affected compared with the gas phase. In this study strong evidence is found for long lived ClO in the electronically

excited  $^2\Pi_{1/2}$  state. A comprehensive IR study of Ne matrix isolated  $\text{ClO}_3$  (fundamentals at 1081, 905, 567,  $476\text{ cm}^{-1}$ ) yield i) a reliable force field; ii) a  $\text{OClO}$  bond angle of  $\alpha_e = 113.8 \pm 1^\circ$  and iii) a  $\text{ClO}$  bond length of  $148.5 \pm 2\text{ pm}$  in agreement with predicted data from quantum chemical calculations. The UV/Vis spectrum of  $\text{ClO}_3$  isolated in a Ne matrix ( $\lambda_{\text{max}}$  at 32100 and  $23150\text{ cm}^{-1}$ ) agrees well with the photoelectron spectrum of  $\text{ClO}_3^-$  and theoretical predictions. The origin of the structured high energy absorption is at

$22696\text{ cm}^{-1}$  and three fundamentals ( $794, 498, 280\text{ cm}^{-1}$ ) are detected in the  $\text{C}^2\text{E}$  state. By photolysis of  $\text{ClO}_3$  with visible light the complex  $\text{ClO}\cdot\text{O}_2$  with ClO in the  $^2\Pi_{1/2}$  state is formed. In an extended spectroscopic study of the elusive  $\text{ClO}_4$  radical, isolated in a Ne or  $\text{O}_2$  matrix, three additional IR bands, a complete UV spectrum and a strong interaction with  $\text{O}_2$  are found. This leads to the conclusion that  $\text{ClO}_4$  exhibits  $\text{C}_{2v}$  or  $\text{C}_s$  symmetry with a shallow potential minimum and forms with  $\text{O}_2$  the previously unknown peroxy radical  $\text{O}_3\text{ClO}-\text{O}_2$ . All these results are discussed in the context of recent developments in the chemistry and spectroscopy of the important and interesting  $\text{ClO}_x$  ( $x = 1-4$ ) family of radicals.

**Keywords:** chlorine compounds · matrix isolation · radical reactions · UV/Vis spectroscopy · vibrational spectroscopy

## Introduction

Binary chlorine oxides  $\text{Cl}_x\text{O}_y$  ( $x = 1, 2$ ;  $y = 1-7$ ) are very important chemical species since they are involved in many reactions of industrial, environmental or academic interest. Due to the environmental role of the mononuclear radicals  $\text{ClO}_x$  ( $x = 1-4$ ), especially in their involvement in ozone-depletion mechanisms and the formation of reactive intermediates, knowledge of these radicals has increased rapidly in recent years.<sup>[1]</sup>

The principal objective of this study is to review the recent developments in the chemistry and spectroscopy of the important and interesting  $\text{ClO}_x$  ( $x = 1-4$ ) family of radicals and to report on new experimental results to complement our previous matrix studies on  $\text{ClO}_3$ <sup>[2]</sup> and  $\text{ClO}_4$ <sup>[3]</sup> radicals and the dimers of  $\text{ClO}$ <sup>[4]</sup> and  $\text{OClO}$ .<sup>[5]</sup>

Within the family of  $\text{ClO}_x$  radicals ( $x = 1-4$ ), chlorine dioxide is the most established species, because it has been known for more than 180 years.<sup>[6-10]</sup> Chlorine dioxide is the only metastable chlorine oxide radical under normal conditions and is also of commercial interest. Its main uses are in the food industries, pulp bleaching in the paper industry, and water purification.<sup>[11]</sup> In the latter case, smaller amounts of chlorinated organic products are formed than by water treatment with chlorine or hypochlorite. There are several routes for the synthesis of  $\text{OClO}$  that involve the reduction of chlorates or the oxidation of chlorites [Eqs. (1), (2)].

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Pure chlorine dioxide is potentially explosive. It can be handled safely as a diluted gas<sup>[12]</sup> (partial pressure up to 400 mbar, temperature below 303 K) or in aqueous solutions. The yellow paramagnetic gas (b.p. 11 °C) crystallizes at -59 °C and the yellow solid becomes diamagnetic at -84 °C.<sup>[13]</sup> The weakly bound dimers in the solid state have been structurally characterized,<sup>[13, 14]</sup> studied in the matrix by IR spectroscopy<sup>[5]</sup> and are the subject of theoretical calculations.<sup>[15]</sup> Scrambling of oxygen atoms between the isotopomers Cl<sup>16</sup>O<sub>2</sub> and Cl<sup>18</sup>O<sub>2</sub> occurs in the liquid state, presumably by dimer formation.<sup>[16]</sup>

For more than 15 years OCIO has become an environmental product which can be observed in the stratosphere over the polar regions during spring time, where it is formed for example by the reaction of ClO with BrO.<sup>[1]</sup> This unexpected occurrence has gained further interest in OCIO. Its properties have been studied extensively, including the microwave,<sup>[17, 18]</sup> millimeter and submillimeter wave,<sup>[19]</sup> high-resolution FTIR,<sup>[20–22]</sup> UV/Vis,<sup>[23, 24]</sup> photoelectron,<sup>[25]</sup> and photoion mass spectra<sup>[26, 27]</sup> and its photochemistry.<sup>[28, 29]</sup> Photofragmentation of OCIO in the gas phase leads to ClO+O or to Cl+O<sub>2</sub>, but the molar ratio of both channels is still under discussion.<sup>[29–35]</sup> Photolysis of matrix isolated OCIO at temperatures below 100 K causes isomerization into ClOO.<sup>[5, 31, 36, 37]</sup> Both endothermic isomers, ClOO and OCIO, are of comparable thermodynamic stability,<sup>[38–41]</sup> but ClOO is a short lived species as a result of its low Cl–O<sub>2</sub> bond energy of only 20 kJ mol<sup>-1</sup>.<sup>[29]</sup> In contrast, organic peroxy radicals R–O<sub>2</sub> are of much higher stability, for example 150 kJ mol<sup>-1</sup> for Et–OO or *t*Bu–OO.<sup>[42]</sup>

Transient ClO radicals have been first identified in studies about the photochemistry of Cl<sub>2</sub>O and OCIO as well as in the thermal and photochemical reactions of ozone with chlorine atoms.<sup>[43–58]</sup> Subsequent introduction of flash-photolysis techniques<sup>[59–61]</sup> has allowed investigation of the UV spectrum of ClO and its kinetics.<sup>[62–70]</sup> It should be noted that Bodenstern and Schumacher studied the kinetics of the reaction between ozone and chlorine atoms first,<sup>[49, 50]</sup> which has become very important in stratospheric chemistry. Ozone depletion processes in the stratosphere induced by chlorine atom release from halocarbons, as claimed by Rowland and Molina in 1974,<sup>[71]</sup> have increased the interest in halogen-oxygen chemistry. For the annual strong ozone depletion in the stratosphere over the polar regions, a catalytic cycle [Eqs. (3)–(5)] has been suggested.<sup>[72]</sup>



Meanwhile the discussion concerning the photochemical cross section of Cl<sub>2</sub>O<sub>2</sub> seems to be finished, since in the cycle mentioned above the contribution of BrOOCl was taken into account.<sup>[73, 74]</sup> Previously found discrepancies in the chlorine budget in the (lower and middle) stratosphere seem to be ruled out by more recent measurements.<sup>[75–79]</sup> Hence, only

sophisticated problems concerning ClO, for example the mechanism of its dimerization, which yields different Cl<sub>2</sub>O<sub>2</sub> isomers,<sup>[4]</sup> the rate constant of the HO<sub>2</sub> + ClO reaction<sup>[80]</sup> or highest rotational levels<sup>[81]</sup> are the subject of further studies. Most synthetic, kinetic, spectroscopic and photochemical aspects of the ClO<sub>*x*</sub> (*x* = 0–2) family are presented in a recent comprehensive review.<sup>[1]</sup>

In the last decades less attention was focussed on the ClO<sub>*x*</sub> (*x* = 3, 4) radicals, although the symmetric ClO<sub>3</sub> radical (C<sub>3v</sub>) has a long history. ClO<sub>3</sub> was postulated as an intermediate in the formation of Cl<sub>2</sub>O<sub>6</sub> from the reaction of Cl or OCIO with O<sub>3</sub>.<sup>[50, 82–85]</sup> and in the decomposition of Cl<sub>2</sub>O<sub>6</sub>.<sup>[86]</sup> The structural, spectroscopic and chemical properties of ClO<sub>3</sub> and Cl<sub>2</sub>O<sub>6</sub> have remained controversial over nearly 150 years.<sup>[47, 87–98]</sup> Finally, for Cl<sub>2</sub>O<sub>6</sub>, a X-ray crystallographic analysis of the low temperature solid has confirmed an ionic structure ClO<sub>2</sub><sup>+</sup>ClO<sub>4</sub><sup>-</sup>,<sup>[99, 100]</sup> while an IR gas phase and matrix isolation study have suggested a molecular structure of O<sub>2</sub>ClOClO<sub>3</sub>.<sup>[101–103]</sup> In the gas phase Cl<sub>2</sub>O<sub>6</sub> decomposes rapidly into OCIO and O<sub>2</sub> and is not in equilibrium with the anticipated ClO<sub>3</sub> radicals.<sup>[103]</sup> Countless attempts to study the formations and spectroscopic characteristics of ClO<sub>3</sub> have been failed. For example the thermal decay of Cl<sub>2</sub>O<sub>6</sub> or Cl<sub>2</sub>O<sub>7</sub> in the presence of fluorine,<sup>[104, 105]</sup> the reaction of ClO with singlet O<sub>2</sub>,<sup>[106]</sup> the reaction between OCIO and O(<sup>3</sup>P),<sup>[28, 68, 107–110]</sup> or the reaction of ClO<sub>*x*</sub> (*x* = 0–2) with O<sub>3</sub>.<sup>[111–114]</sup>

In 1994 ClO<sub>3</sub> radicals have been prepared for the first time by the low pressure pyrolysis of ClOClO<sub>3</sub> [Eq. (6)] followed by subsequent quenching of the products in low temperature matrices and are unambiguously characterized by IR and UV spectroscopy.<sup>[2]</sup>



ClO<sub>3</sub> shows a broad and structured UV/Vis spectrum beginning at about 750 nm with increasing intensity towards 300 nm. All four expected fundamentals for the <sup>2</sup>A<sub>1</sub> ground state have been measured including the <sup>35</sup>S/<sup>37</sup>Cl and <sup>16</sup>O/<sup>18</sup>O isotopic species. Finally the formation of ClO<sub>3</sub> was supported by the excellent agreement between observed and calculated ab initio IR band positions and intensities.<sup>[2, 38–40, 115–119]</sup> In addition, time resolved UV spectra of pulse and  $\gamma$ -radiolysed or laser photolysed chlorate or perchloric acid solutions are assigned to ClO<sub>3</sub> or ClO<sub>4</sub> radicals, respectively.<sup>[120–122]</sup>

The existence of ClO<sub>4</sub> has also been controversial in the last 80 years as discussed;<sup>[85, 104, 105, 111, 113, 123–132]</sup> it was only in the late eighties of the last century that reinvestigations on the decay of Cl<sub>2</sub>O<sub>6</sub><sup>[101, 103]</sup> lead to the trustworthy assumption that an intermediate ClO<sub>4</sub> radical is formed [Eq. (7.1)].



The first IR spectra of matrix isolated ClO<sub>4</sub> have been reported in 1996.<sup>[3]</sup> The ClO<sub>4</sub> radical was prepared by low pressure thermolysis of Cl<sub>2</sub>O<sub>6</sub> with subsequent quenching of the products at low temperature in matrix. From the limited spectral data for ClO<sub>4</sub>, C<sub>3v</sub> symmetry was deduced. However, several new theoretical calculations on highest level predict C<sub>2v</sub> or C<sub>s</sub> symmetry.<sup>[39, 40, 119, 133, 134]</sup> Recently kinetic measure-

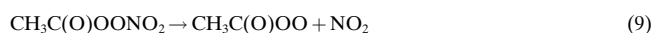
ments of the decay of  $\text{Cl}_2\text{O}_4$  by reaction with chlorine atoms [Eq. (8)] lead to the assumption, that a stationary partial pressure of about 0.07 mbar  $\text{ClO}_4$  may be possible in a steady state system.<sup>[135]</sup>



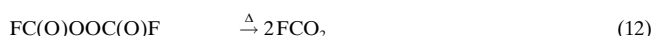
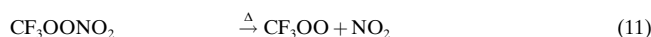
Another possible access to  $\text{ClO}_4$  by  $\gamma$ -irradiation of solid perchlorates is promising, but the discussion and interpretation of the detected ESR signals have been controversial.<sup>[96, 97, 136–151]</sup>

## Results and Discussion

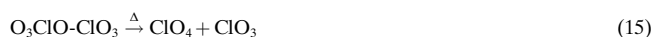
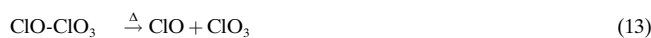
The spectroscopic properties of short-lived radicals or molecules can be easily investigated if they are isolated in cryogenic matrices.<sup>[152–154]</sup> In a matrix-isolation experiment a high vacuum and the use of highly diluted precursor molecules during matrix deposition offer the appropriate conditions for the synthesis of short-lived radicals and molecules by low pressure flash thermolysis. In low pressure flash thermolysis an unimolecular bond fission of a fragile molecule A–B can be easily achieved, if the A–B bond energy is below  $200 \text{ kJ mol}^{-1}$ . If the bonds within both fragments A and B are stronger than  $150 \text{ kJ mol}^{-1}$ , little secondary bond cleavage of the A or B radicals should occur. These conditions are found to be true for several typical examples, that we studied in recent years. An early example is the fragmentation of peroxyacetyl nitrate (PAN)<sup>[155]</sup> [Eq. (9)] which leads to new spectroscopic data of the peroxyacetyl radical.



Recent prominent experiments are the formation of  $\text{CF}_3\text{O}$ ,  $\text{CF}_3\text{OO}$  and  $\text{FCO}_2$  radicals by cleaving the thermally fragile precursors  $\text{CF}_3\text{OC}(\text{O})\text{OOC}(\text{O})\text{OCF}_3$ ,<sup>[156]</sup>  $\text{CF}_3\text{OONO}_2$ <sup>[157, 158]</sup> and  $\text{FC}(\text{O})\text{OOC}(\text{O})\text{F}$ <sup>[159]</sup> with subsequent matrix isolation of the products according to Equations (10)–(12).



In order to generate  $\text{ClO}$ ,  $\text{ClO}_3$  or  $\text{ClO}_4$  radicals in a similar manner, suitable precursor molecules containing a  $\text{ClO}$ ,  $\text{ClO}_3$  or  $\text{ClO}_4$  fragment are needed. For this purpose we have investigated the low pressure flash thermolysis of  $\text{XOClO}_3$  ( $\text{X} = \text{H}, \text{F}, \text{Cl}, \text{ClO}_2, \text{ClO}_3, \text{NO}_2$ ) and  $\text{ClONO}_2$  molecules. Other possible precursors  $\text{YOCIO}_2$  ( $\text{Y} = \text{H}, \text{F}, \text{Cl}, \text{ClO}_2, \text{NO}_2$ ) are either unknown or extremely unstable and, therefore not available. In X–O– $\text{ClO}_3$  molecules, the terminal Cl–O bonds in  $\text{ClO}_3$  are very strong and bond cleavage is possible only between  $\text{XO–ClO}_3$  or  $\text{X–OCIO}_3$  [Eq. (13)–(15)].



Which bond is cleaved, depends on the stability of the formed fragments. Therefore, in  $\text{ClOClO}_3$  or  $\text{O}_2\text{ClOClO}_3$  the stable fragment radicals  $\text{ClO}$  or  $\text{ClO}_2$  are expected to be formed and the  $\text{ClO}_3$  or  $\text{ClO}_4$  radical should be the counterpart, respectively.

The thermal dissociation can be predicted from the molecular structures and bond energies. The structures of  $\text{XOClO}_3$  molecules have been investigated by electron diffraction in the gas phase. In all cases the longest bond for  $\text{XO–ClO}_3$  ( $\text{X} = \text{H}, \text{F}, \text{Cl}$ ) amounts 164, 170, 171 pm, respectively<sup>[160]</sup> and the related bond energies are estimated to be 250, 181 and  $150 \text{ kJ mol}^{-1}$ , respectively.<sup>[161, 162]</sup> Therefore, the best precursor for the  $\text{ClO}_3$  radical should be  $\text{ClO–ClO}_3$ .

Indeed, the best results have been achieved by the thermolysis of  $\text{ClOClO}_3$  or  $\text{FOClO}_3$ . In regard to by-products we have used,  $\text{ClOClO}_3$  for recording IR spectra and  $\text{FOClO}_3$  for recording the UV spectra, because the UV absorption of the  $\text{XO}$ ,  $\text{XOO}$  or  $(\text{XO})_2$  by-products for  $\text{X} = \text{F}$  are blue-shifted in comparison to the species with  $\text{X} = \text{Cl}$  and they interfere less with the UV absorptions of  $\text{ClO}_3$ . At higher temperatures, which are needed for the cleavage of  $\text{HO–ClO}_3$  or  $\text{O}_3\text{Cl–OCIO}_3$ , increased amounts of  $\text{ClO}_2$  are formed.

Covalently bound  $\text{Cl}_2\text{O}_6$  isomerizes into ionic  $\text{ClO}_2^+\text{ClO}_4^-$  in the solid state below  $-5^\circ\text{C}$ ,<sup>[101, 102]</sup> the related  $\text{NO}_2\text{ClO}_4$  is even ionic ( $\text{NO}_2^+\text{ClO}_4^-$ ) at room temperature.<sup>[163]</sup> As discussed above such  $\text{X–OCIO}_3$  ( $\text{X} = \text{ClO}_2, \text{NO}_2$ ) molecules are promising precursors for the  $\text{ClO}_4$  radical. The molecular structure of  $\text{O}_3\text{Cl–OCIO}_3$  with an O–Cl bond length of 172 pm indicates a weak bond<sup>[164, 165]</sup> and hence it may be a precursor for both  $\text{ClO}_3$  and  $\text{ClO}_4$  radicals. The estimated bond energies for  $\text{O}_2\text{Cl–OCIO}_3$ ,  $\text{O}_3\text{Cl–OCIO}_3$ , and  $\text{O}_2\text{N–OCIO}_3$  are 100, 140 and  $92 \text{ kJ mol}^{-1}$ , respectively.<sup>[161, 162]</sup> Therefore the best precursor for  $\text{ClO}_4$  radicals should be  $\text{NO}_2\text{ClO}_4$ . However, the vapour pressure of  $\text{NO}_2\text{ClO}_4$  at room temperature is not sufficient for matrix isolation and at elevated temperatures solid  $\text{NO}_2\text{ClO}_4$  decomposes. Another possible precursor for  $\text{ClO}_4$  may be  $\text{O}_3\text{ClO–OCIO}_3$ , because the isoelectronic  $\text{FSO}_2\text{O–OSO}_2\text{F}$  is well known and an excellent source for  $\text{FSO}_3$  radicals.<sup>[166]</sup> However, all attempts for the synthesis of  $\text{O}_3\text{ClOOCIO}_3$  or the mixed peroxide  $\text{FSO}_2\text{OOCIO}_3$  have failed so far. The best precursor for  $\text{ClO}$  radicals is  $\text{ClONO}_2$  due to its low  $\text{ClO–NO}_2$  bond energy of  $109 \text{ kJ mol}^{-1}$ .<sup>[167]</sup>

Photoreactions in matrices are a powerful tool to create new species, which are not accessible by other reaction pathways. The cage effect is one of the most important features of matrix isolation. This effect is not available in the gas phase. The central question is, whether the photo fragments escape from the matrix cage and if not, how do the fragments rearrange? If a fragment is small and symmetrical, for example a hydrogen atom ( $^2\text{S}_{1/2}$ ), it may escape from the cage. In contrast, halogen ( $^2\text{P}_{3/2}$ ) or oxygen ( $^3\text{P}_{3/2}$ ) atoms are nonsymmetrical in their ground state. Hence, they should hardly escape but add again to a favoured side of the parent fragment; this results in an isomerization of the parent species. The size of the cage is the limiting factor of this photoinduced process.

Several photochemical isomerization reactions of chlorine oxides are well known: The monomer of chlorine dioxide  $\text{OCIO}$  transforms into the chlorine superoxide  $\text{ClOO}$  on irradiation with light of wavelength  $\leq 475 \text{ nm}$ .<sup>[31]</sup> In this case

the O atom is the mobile agent, linking to the O atom of the parent fragment. An example for a mobilised Cl atom is the Y-shaped chloryl chloride,  $\text{ClClO}_2$ . It shows a typical cage effect: In the small cage of a neon matrix  $\text{Cl}_2$  and  $\text{O}_2$  are formed on irradiation at  $\lambda \geq 495$  nm, while in argon matrix there is enough space for a rotational motion of the  $\text{OCIO}$  fragment. At  $\lambda \leq 610$  nm  $\text{ClClO}_2$  isomerises into chlorine chlorite  $\text{ClOClO}$ , which undergoes a further transformation into dichlorine peroxide  $\text{ClOClO}$ .<sup>[4]</sup> Johnsson et al. have shown that this is also true for  $\text{XClO}_2$  with  $\text{X} = \text{Br}$  while for  $\text{X} = \text{I}$  even the argon cage is too small.<sup>[168]</sup>

For a deeper understanding of the photo isomerisation mechanism, we have reinvestigated the photolysis of  $\text{ClO}_x$  ( $x = 1-4$ ) in oxygen matrices. The advantage of these experiments is that the mobile O atom will be caught mainly by the reactive oxygen matrix material forming ozone. A fragment is left behind, to which the O atom would have been added in the case of a non-reactive matrix.

### Chlorine Monoxide

Many spectroscopic studies on gaseous  $\text{ClO}$ , which is a weak IR<sup>[1, 169, 170]</sup> and strong UV<sup>[1, 68-70]</sup> absorber, have been performed in the past. We want to report here on the influence of different matrix environments on the energetic states of  $\text{ClO}$  and on recent observations of its dimerization. In addition, the recorded spectra will be used as reference in the analysis of the spectra of  $\text{ClO}_3$  and  $\text{ClO}_4$  discussed below.

Low pressure flash pyrolysis of  $\text{ClONO}_2$ , highly diluted in a matrix gas with subsequent quenching of the products on a cold matrix support, leads to  $\text{ClO}$ ,  $\text{NO}_2$  and minor amounts of  $\text{ClOClO}$  according to the IR spectra. Hence, the thermolysis may proceed as follows [Eq. (16)–(19)].



secondary:



or:



**The IR spectrum of  $\text{ClO}$ :** The vibrational wavenumber of  $^{35}\text{ClO}$  is found to show very little dependency on the environment (the respective  $^{37}\text{ClO}$  band position is in all cases identical to the calculated value from the two-mass model). In Ne matrix ( $843.8 \text{ cm}^{-1}$ ) it is  $0.5 \text{ cm}^{-1}$  red-shifted in comparison to the gas-phase value of  $844.20 \text{ cm}^{-1}$ <sup>[171, 172]</sup> and in other matrix materials (Ar:  $844.9$ ;  $\text{N}_2$ :  $848.2$ ; Kr:  $848.8$ ;  $\text{O}_2$ :  $849.4 \text{ cm}^{-1}$ ) blue-shifts are observed. This trend is unusual, because there are very often increased red-shifts of stretching modes of matrix isolated molecules with increasing polarisability of the matrix material.<sup>[173]</sup> Because the oxygen matrix behaves like an inert matrix material, there should be only a weak chemical interaction between  $\text{ClO}$  and oxygen. Con-

cerning the interaction between  $\text{ClO}$  and  $\text{O}_2$ , interesting observations summarized in our previous  $\text{ClO}_3$  paper<sup>[2]</sup> are now confirmed and extended by new and improved experiments. By UV photolysis of Ne matrix isolated  $\text{ClO}_3$  radicals two new  $^{35}\text{ClO}$  bands at  $844.0$  and at  $839.5 \text{ cm}^{-1}$  appear; the latter is tentatively assigned to a  $\text{ClO} \cdot \text{O}_2$  complex. However, by comparing with the  $^{35}\text{ClO}$  wavenumber in oxygen matrix ( $849.4 \text{ cm}^{-1}$ ) one would expect a significant blue shift for a  $\text{ClO} \cdot \text{O}_2$  complex. Since the former band is only slightly blue-shifted ( $+0.2 \text{ cm}^{-1}$ ), this band is assigned to  $\text{ClO}$  radicals, isolated in a Ne matrix cage, which are little disturbed by an oxygen molecule. The latter band shows a red shift ( $-4.3 \text{ cm}^{-1}$ ) and is assigned to a  $\text{ClO} \cdot \text{O}_2$  complex. For further details on this puzzle see section below on chlorine trioxide. There is another interesting feature in the IR spectra of  $\text{ClO}$  isolated in Ne matrices from the gas phase (without any  $\text{O}_2$ ): A weak additional band at  $838.5 \text{ cm}^{-1}$  is always observed, which is close to the band of the  $\text{ClO} \cdot \text{O}_2$  complex at  $839.5 \text{ cm}^{-1}$ . A possible interpretation is that this band is due to  $\text{ClO}$  in the  $^2\Pi_{1/2}$  state. This assignment arises from the analysis of the UV spectrum of Ne matrix isolated  $\text{ClO}$  (see below) where a  $\text{A}^2\Pi_{1/2} \leftarrow \text{X}^2\Pi_{1/2}$  subsystem is observed. This is only possible, if the  $\text{X}^2\Pi_{1/2}$  state, which is higher in energy by  $318 \text{ cm}^{-1}$  than the  $\text{X}^2\Pi_{3/2}$  ground state,<sup>[174]</sup> is still present after matrix deposition. Therefore we assume for the  $\text{X}^2\Pi_{1/2}$  state of matrix isolated  $\text{ClO}$  a relaxation time of several hours. In the gas phase the vibrational wavenumber of the  $\text{X}^2\Pi_{1/2}$  state is about  $3.3 \text{ cm}^{-1}$  lower than that of the  $\text{X}^2\Pi_{3/2}$  state.<sup>[171, 172]</sup> In the Ne matrix IR spectra of  $\text{ClO}$ , produced by pyrolysis of  $\text{ClONO}_2$  or  $\text{ClOClO}_3$ , a weak satellite  $5.3 \text{ cm}^{-1}$  below the absorption of  $^{35}\text{ClO}$  in the  $\text{X}^2\Pi_{3/2}$  state indicates the presence of a small fraction of  $^{35}\text{ClO}$  in the  $\text{X}^2\Pi_{1/2}$  state.

**The UV spectrum of  $\text{ClO}$ :** The UV spectrum of the pyrolysis products of  $\text{ClONO}_2$ , isolated in a neon matrix, is shown in Figure 1. From this spectrum, reference spectra of  $\text{NO}_2$  and residual  $\text{ClONO}_2$  have been subtracted, until a proper baseline for the final spectrum of  $\text{ClO}$  results. An identical UV spectrum can be obtained in a similar manner from the precursor  $\text{ClOClO}_3$ . The cross section scale in Figure 1 is adopted from the  $\text{NO}_2$  gas phase values<sup>[157, 158]</sup> and the

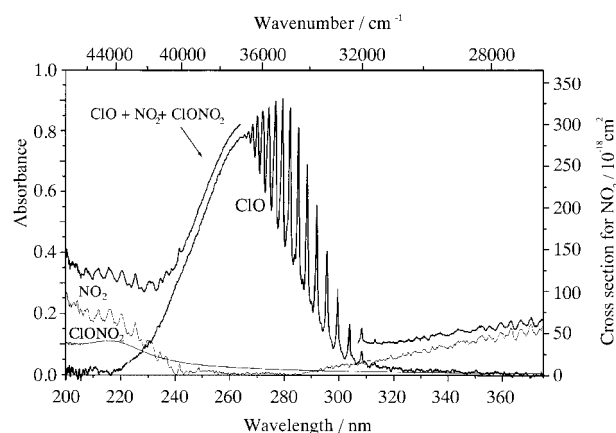


Figure 1. UV spectrum of  $\text{ClO}$  isolated in a neon matrix. Upper trace: raw spectrum of the thermolysis product of  $\text{ClONO}_2$ . Middle trace: reference spectrum of  $\text{NO}_2$  and  $\text{ClONO}_2$  (lower trace).

resulting absorption cross sections of ClO are listed in Table 1. Assuming that ClO and NO<sub>2</sub> are produced at a 1:1 molar ratio according to Equation (16) the cross sections for ClO are estimated to be too low. The difference in the cross section of matrix isolated ClO compared with the gas-phase values (Table 1) is attributed to side reactions during the thermolysis and quenching process [Eqs. (17)–(19)]. Due to the low concentration of ClOClO and other undetectable ClO dimers, these species should not contribute to the raw UV spectrum. At the absorbance maximum of Cl<sub>2</sub> near 330 nm no interference of the ClO band profile can be seen in Figure 1. Hence the main loss of ClO must occur in the heated nozzle with formation of chlorine atoms and dioxygen [Eq. (19)].

In the gas phase, the UV cross sections of the continuum is independent of temperature, while the structured part is extremely resolution and temperature dependent.<sup>[1, 69]</sup> Comparing our NO<sub>2</sub> based Ne matrix cross sections with literature gas-phase data at the beginning of the continuum at 260 nm, these values can be adjusted with a scaling factor of 1.88 (Table 1). Then the resulting values of the continuum are in excellent agreement with literature data,<sup>[62, 66, 69]</sup> whereas the values of Simon et al.<sup>[70]</sup> show a slower decrease at shorter wavelengths. In comparison to the structured part of the gas-phase spectrum, there are no hot band transitions and no rotational fine structures in the Ne matrix spectrum. Very surprising is the presence of the A <sup>2</sup>Π<sub>1/2</sub> ← X<sup>2</sup>Π<sub>1/2</sub> subsystem, as already mentioned above. The scaled cross sections of the progression bands are in good agreement with literature values (see Table 1) up to the vibrational excited state (9,0).

Table 1. Absorption cross sections of ClO (A<sup>2</sup>Π<sub>3/2</sub> ← X<sup>2</sup>Π<sub>3/2</sub>) isolated in a Ne matrix and in the gas phase.

Transition	Ne matrix				Gas phase			[f]
	λ <sup>[a]</sup> /(nm)	σ <sup>[b]</sup> /(10 <sup>-20</sup> cm <sup>2</sup> ) slit 0.3 nm		λ/(nm)	σ/(10 <sup>-20</sup> cm <sup>2</sup> )			
					0.3 nm <sup>[c]</sup>	0.3 nm <sup>[d]</sup>	0.22 nm <sup>[e]</sup>	
1,0	313.2 (312)	15.1	28.4	312.5			26	
2,0	308.4 (305.7)	30.5	57.3	307.9	39	58	47	
3,0	303.9 (301.6)	59.3	112	303.5	86	103	104	
4,0	299.7 (297.5)	103	194	299.3	163	171	104	
5,0	295.8 (293.8)	153	288	295.4	255	253	207	
6,0	292.0 (290.2)	205	385	292.0	338	330	382	
7,0	288.5 (287.0)	257	483	288.4	448	438	516	
8,0	285.3 (283.9)	296	556	285.2	542	530	627	
9,0	282.3 (281.1)	322	605	282.3	608	598	667	
10,0	279.4 (278.5)	332	624	279.8	655	645	688	
11,0	276.9 (276.1)	330	620	277.2	679	675	680	
12,0	274.5 (274.0)	325	611	275.1	681	671	634	
13,0	272.2 (272)	320	602					
14,0	270.2 (271)	314	590					
15,0	268.5 (269)	304	572					
16,0	266.9 (267)	295	555					
17,0	265.5 (266)	291	547					
continuum	260.0	270	508	260.0	504	509	512	530
	257.7	256	481	257.7	488	480	485	485
	255.0	235	442	255.0	454	440	452	450
	250.0	191	359	250.0	376	351	352	360
	245.0	145	273	245.0	294	274		270
	240.0	103	194	240.0	222			190
	235.0	68	128					
	230.0	42	79					
	225.0	28	53					
	220.0	16	30					

[a] Values in parenthesis belong to the A<sup>2</sup>Π<sub>1/2</sub> ← X<sup>2</sup>Π<sub>1/2</sub> subsystem. [b] Second values scaled with 1.88, see text. [c] Ref. [70]. [d] Ref. [69]. [e] Ref. [66]. [f] Ref. [65].

Due to interactions with the matrix phonon bands, the bands become broader and less intensive at higher excitations. Another reason for the band broadening is the transitions of the A<sup>2</sup>Π<sub>1/2</sub> ← X<sup>2</sup>Π<sub>1/2</sub> subsystem, which interfere at higher excitations with the A<sup>2</sup>Π<sub>3/2</sub> ← X<sup>2</sup>Π<sub>3/2</sub> system.

The progression of both subsystems has been analysed and the resulting spectroscopic data of the excited states are compared with the gas-phase data in Table 2. Both, the band

Table 2. Spectroscopic data [cm<sup>-1</sup>] of <sup>35</sup>Cl<sup>16</sup>O in the gas phase and isolated in a Ne matrix.

Ground state	Excited state	ν (0,0)	ω <sub>e</sub>	χ <sub>e</sub> ω <sub>e</sub>
gas phase				
X <sup>2</sup> Π <sub>3/2</sub>	A <sup>2</sup> Π <sub>3/2</sub> ← X <sup>2</sup> Π <sub>3/2</sub>	31 482 <sup>[b]</sup>	520 <sup>[b]</sup>	7.2 <sup>[b]</sup>
X <sup>2</sup> Π <sub>1/2</sub>	A <sup>2</sup> Π <sub>1/2</sub> ← X <sup>2</sup> Π <sub>1/2</sub>	31 685 <sup>[b]</sup>		
Ne matrix <sup>[c]</sup>				
X <sup>2</sup> Π <sub>3/2</sub>	A <sup>2</sup> Π <sub>3/2</sub> ← X <sup>2</sup> Π <sub>3/2</sub>	31 408	526	8.9
X <sup>2</sup> Π <sub>1/2</sub>	A <sup>2</sup> Π <sub>1/2</sub> ← X <sup>2</sup> Π <sub>1/2</sub>	31 670	527	10.3

[a] Ref. [171, 172]. [b] Ref. [174]. [c] This work.

origin and the vibrational wavenumber of ClO are only slightly affected by the matrix, while the anharmonicity is strongly increased. The latter effect is also seen in Table 1: The maxima of the progression bands are red-shifted at the beginning and blue-shifted at higher quantum numbers. This trend can be rationalized by increased repulsion of the ClO motions on the matrix with higher quantum numbers.

**Dimerization of ClO:** As mentioned above, some dimerization of ClO is always observed during deposition of ClO in a matrix. Chainlike ClOClO is formed by head-to-tail interaction of two ClO and the resulting ClOClO shows a characteristic, strong IR band at 994.5 and 985.8 cm<sup>-1</sup> (<sup>35/37</sup>ClO<sup>35</sup>ClO, <sup>35/37</sup>ClO<sup>37</sup>ClO, Ar matrix). This dimer has been observed already in very early matrix isolation experiments of ClO.<sup>[175–178]</sup> In later quantum chemical calculations it was pointed out, that three different isomers ClOCl, ClClO<sub>2</sub> and ClOClO should exist with decreasing thermodynamic stability, respectively,<sup>[179–182]</sup> of these, ClClO<sub>2</sub> has the highest kinetic stability, which can be prepared by F/Cl exchange from FClO<sub>2</sub>.<sup>[183, 184]</sup> Matrix-isolated ClClO<sub>2</sub> can be converted photolytically into the two other isomers.<sup>[4, 184]</sup> In the homogenous gas phase, ClOCl is formed by

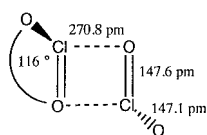
dimerization of ClO and characterized by microwave spectroscopy.<sup>[185]</sup> On the other hand for example on ice surfaces ClO dimerization leads to the isomer ClClO<sub>2</sub>.<sup>[186]</sup> Hence, just which species of the three Cl<sub>2</sub>O<sub>2</sub> isomers are formed, strongly depends on the reaction conditions.

## Chlorine Dioxide

Since the first detailed IR matrix isolation study on OCIO and its photolysis product ClOO,<sup>[187]</sup> many further spectroscopic investigations on both isomers in different matrices have been performed. IR absorption spectra of OCIO and its isomer ClOO, isolated in inert matrices<sup>[5, 188]</sup> and in crystalline ice,<sup>[36, 37, 189]</sup> have been studied extensively. The UV spectra of ClOO<sup>[5, 188]</sup> and of OCIO<sup>[5, 31]</sup> isolated in a Ne matrix have been measured and analysed. Recently, dispersed laser-induced fluorescence spectra of the A<sup>2</sup>A<sub>2</sub> → X<sup>2</sup>B<sub>1</sub> transition of OCIO in solid Ne have been recorded.<sup>[190]</sup> Mode specificity by photodissociation in a molecular beam<sup>[29]</sup> and site selectivity by laser photolysis of OCIO in solid rare gases have been observed and the influence of different matrices on the conversion factors  $\Delta A_{\text{ClOO}}/(-\Delta A_{\text{OCIO}})$  were discussed.<sup>[31]</sup>

In all these studies appreciable amounts of OCIO dimers are present in the matrices, which caused some problems in the interpretation of the results. In order to contribute to the understanding of the (OCIO)<sub>2</sub> interference, we want to report here mainly on the (OCIO)<sub>2</sub> dimer and in addition on the photochemistry of ClO<sub>2</sub> and (OCIO)<sub>2</sub>, isolated in an oxygen matrix.

**The IR spectrum of the (OCIO)<sub>2</sub> dimer:** In the solid state, OCIO forms weakly bound dimers as depicted of C<sub>i</sub> symmetry with the following structural parameters.<sup>[13]</sup> These dimers partially survive sublimation at -120 °C and can be trapped in



a matrix within a few milliseconds. We believe, that they maintain a similar structure. In this manner it has been possible to prepare matrices with a high fraction on OCIO dimers. By heating the spray on nozzle in front of the matrix support and by using more highly diluted matrices, the contents of dimers are strongly reduced. Typical IR spectra for such an experiment are depicted in Figure 2.

The 12 vibrational motions of the OCIO dimer are subdivided according to the irreducible representation [Eq. (20)].

$$\Gamma_{\text{vib}} = 6A_g (\text{Ra, p}) + 6A_u (\text{IR}) \quad (20)$$

The two IR active ClO stretching modes and the OCIO bending mode are observed in close proximity to the fundamentals of the monomer. All other fundamentals are expected to be observed in the far IR region and in the Raman spectrum. The band at 1107 cm<sup>-1</sup> is obscured by the fundamental of the monomer and at 935 cm<sup>-1</sup>, the 9:6:1 <sup>35/37</sup>Cl isotopic pattern is typical for a species, containing two equivalent chlorine atoms. The experimental IR data of (OCIO)<sub>2</sub> are collected in Table 3.

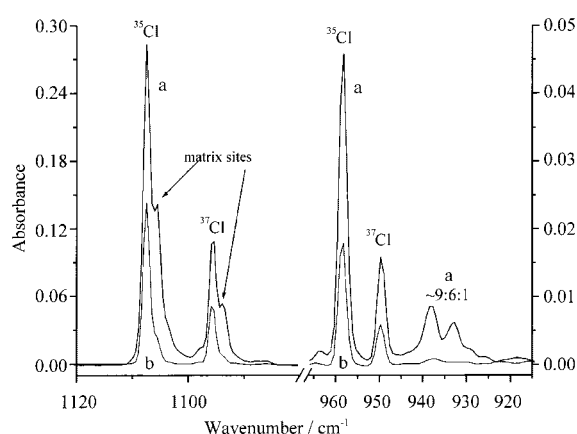


Figure 2. IR spectrum of OCIO isolated in a neon matrix. a) OCIO with a high content of (ClO<sub>2</sub>)<sub>2</sub> dimers. b) Monomeric OCIO.

Table 3. Experimental vibrational wavenumbers [cm<sup>-1</sup>] for (ClO<sub>2</sub>)<sub>2</sub>.

Mode	Crystalline 77 K and 193 K <sup>[a]</sup>	Neon matrix <sup>[b]</sup>
$\nu_{\text{as}} \text{ClO}_2$	1062 vbr.	1063 vbr
$\nu_{\text{s}} \text{ClO}_2$	908	913
$\delta_{\text{s}} \text{ClO}_2$	471	465
$\nu \text{Cl}_2\text{O}_2$	187	177
		1107 1096
		931, 927, 925/ ≈ 9:6:1
		464 br, 451 br/ ≈ 9:6
		175 br

[a] Ref. [13]. [b] Ref. [5] and this work.

**The UV spectrum of the (OCIO)<sub>2</sub> dimer:** Figure 3 presents the UV/Vis spectra of the same samples, which were used for recording the IR spectra in Figure 2. Traces a and b possess the known progression pattern of OCIO, but in trace b an underlying continuum absorption is present. The weighted

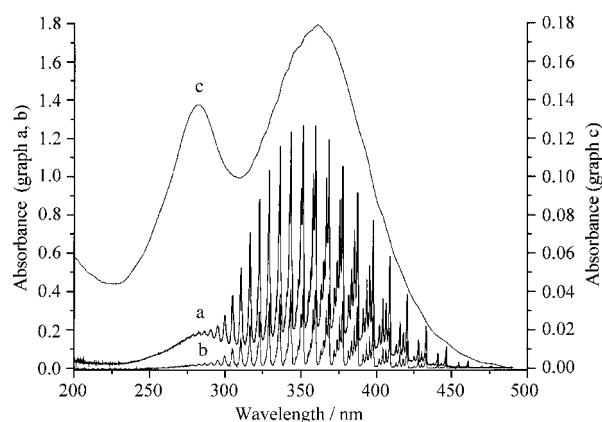


Figure 3. UV/Vis spectrum of OCIO isolated in a neon matrix. a) OCIO with a high content on (ClO<sub>2</sub>)<sub>2</sub> dimers. b) Monomer OCIO. c) Difference spectrum a) - b); (ClO<sub>2</sub>)<sub>2</sub> in expanded scale.

difference b - a which yields the smoothed trace c represents the unstructured UV/Vis spectrum of the (OCIO)<sub>2</sub> dimer. A continuum absorption of OCIO is mentioned by Wahner et al.<sup>[23]</sup> but its origin and intensity are not well understood. Liu et al.<sup>[190]</sup> report recently, that their OCIO continuum might be due to an absorption of the 1<sup>2</sup>A<sub>1</sub> state above its predissociative barrier, but we suppose, that there is a continuum due to dimeric chlorine dioxide.

**Photochemistry of matrix isolated OCIO:** In order to understand the different conversion factors  $\Delta A_{\text{ClOO}}/(-\Delta A_{\text{OCIO}})$  in

dependence of the matrix material and the excitation energy,<sup>[31]</sup> we have studied the OCIO photolysis also in an oxygen matrix. There is no obvious interaction of OCIO and its isomer CIOO with oxygen, because with respect to the wavenumbers of the OCIO and CIOO fundamentals, the oxygen matrix behaves like an argon matrix. However, the photochemistry of OCIO in solid oxygen is quite different from that in solid rare gases. After primary photodissociation of OCIO into CIO+O within the oxygen matrix cage, the surrounding O<sub>2</sub> molecules compete with CIO in the reaction with the oxygen atom. According to the resulting IR spectra the main product is O<sub>3</sub>. In addition some CIOO and the by-product CIO is clearly seen. The ratio of the product band absorbances, O<sub>3</sub>/CIOO, increases with increasing energy of the photolysis light. This means, that the oxygen atom escapes more efficiently from the matrix cage with increasing excess energy of the photolysis light source. This loss of oxygen atoms must be taken in account also in solid rare gases. It becomes more pronounced in the order Ne < Ar < Kr. The (OCIO)<sub>2</sub> dimer behaves similarly. In rare gas matrices it is converted mainly to CIOClO<sub>3</sub> and a small fraction of CIOClO<sub>2</sub>.<sup>[5]</sup> But in an oxygen matrix, the yield of CIOClO<sub>2</sub> is much higher, because of effective oxygen atom scavenge by the surrounding oxygen molecules. UV photolysis ( $\lambda > 230$  nm) of the O<sub>2</sub> matrix for a longer period of time causes a decrease of bands due to CIOClO<sub>3</sub>, CIOClO<sub>2</sub> and CIO and increase of the bands of CIOO. Hence the final end products Cl+O<sub>2</sub> of all chlorine oxides are formed.

### Chlorine Trioxide

In the course of our preliminary study of matrix-isolated CIO<sub>3</sub> radicals, IR and UV/Vis spectra have been measured and analysed to only some extent.<sup>[2]</sup> In particular the reported UV spectrum is strongly perturbed by absorptions due to the by-products CIO and OCIO and to the precursor CIOClO<sub>3</sub>. We want to present here further spectroscopic data and their detailed analysis, in order to characterize this short lived radical more comprehensively.

**The IR spectrum of CIO<sub>3</sub>:** Low pressure flash pyrolysis of CIOClO<sub>3</sub>, highly diluted in a matrix gas, in front of the matrix support with subsequent quenching of the products in matrix leads to CIO, CIO<sub>3</sub>, some CIO<sub>2</sub> and minor amounts of CIOClO, Cl<sub>2</sub>O<sub>6</sub> and Cl<sub>2</sub>O<sub>3</sub>. The formation of these products is assumed to proceed as follows [Eqs. (21)–(26)].

primary:



secondary:



Figure 4 shows a typical IR spectrum of matrix isolated CIO, and CIO<sub>3</sub> as products after subtracting the bands of CIO<sub>2</sub> and residual CIOClO<sub>3</sub> (the ranges are indicated by arrows). All band positions and their relative integrated intensities of CIO<sub>3</sub>

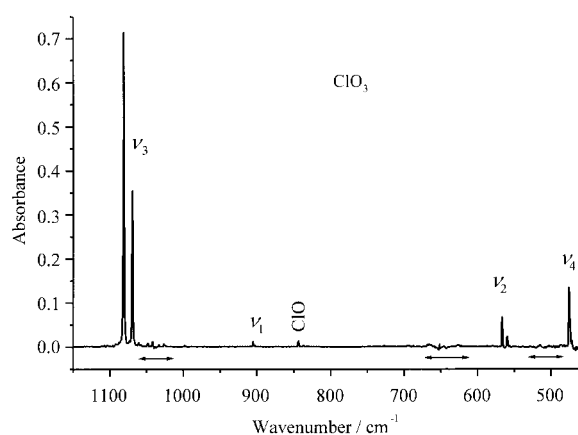


Figure 4. IR spectrum of CIO<sub>3</sub> and CIO isolated in a neon matrix. The radicals are produced by low pressure thermolysis of CIOClO<sub>3</sub>. Bands of residual CIOClO<sub>3</sub> and of some by-products are subtracted in the regions indicated by arrows.

as well as of the Cl<sup>18</sup>O<sub>3</sub> isotopomer are collected in Table 4. Also included are the overtones and combinations, which have not been reported before.<sup>[2]</sup> The band positions are shifted to lower wavenumbers from Ne to Ar or O<sub>2</sub> matrix, as expected for matrices with higher polarizability, and O<sub>2</sub> behaving similar to an Ar matrix. Hence CIO<sub>3</sub> does not react with O<sub>2</sub> to form the peroxy radical O<sub>3</sub>CIOO as does the isoelectronic CF<sub>3</sub> radical,<sup>[157, 191, 192]</sup> both radicals have 25 valence electrons. The identity of all observed CIO<sub>3</sub> bands is confirmed by their constant relative band intensities in

Table 4. Vibrational wavenumbers [cm<sup>-1</sup>] of isotopic CIO<sub>3</sub> species exhibiting C<sub>3v</sub> symmetry isolated in Ne, Ar and O<sub>2</sub> matrix.

<sup>35</sup> Cl <sup>16</sup> O <sub>3</sub>	I <sup>[c]</sup>	Ne matrix <sup>[a]</sup>			Ar matrix <sup>[a]</sup>		O <sub>2</sub> matrix <sup>[a]</sup>		Assignment <sup>[b]</sup>
		<sup>37</sup> Cl <sup>16</sup> O <sub>3</sub>	<sup>35</sup> Cl <sup>18</sup> O <sub>3</sub>	<sup>37</sup> Cl <sup>18</sup> O <sub>3</sub>	<sup>35</sup> Cl <sup>16</sup> O <sub>3</sub>	<sup>37</sup> Cl <sup>16</sup> O <sub>3</sub>	<sup>35</sup> Cl <sup>16</sup> O <sub>3</sub>	<sup>37</sup> Cl <sup>16</sup> O <sub>3</sub>	
2155.31	3.5	2131.77	2087.0	2062.4	2137.0	2113.5	2139.8	2116.2	2ν <sub>3</sub> , A <sub>1</sub>
1970.4	13	1956.3	1892.7	1877.7	1955.3	1941.2	1958.3	1944.1	ν <sub>1</sub> + ν <sub>3</sub> , E
1641.9	0.7	1623.0	1590.7	1571.3	1640.4	1620.8	1630.8	1612.0	ν <sub>2</sub> + ν <sub>3</sub> , E
1081.27	100	1069.44	1047.06	1034.69	1072.2	1060.3	1074.0	1062.2	ν <sub>3</sub> , e ν <sub>as</sub> (ClO <sub>3</sub> )
905.04	3	902.59	860.23	857.45	899.5	897.0	900.4	897.8	ν <sub>1</sub> , a <sub>1</sub> ν <sub>s</sub> (ClO <sub>3</sub> )
566.63	11	559.44	549.63	542.48	562.7	555.6	563.0	555.9	ν <sub>2</sub> , a <sub>1</sub> δ <sub>s</sub> (ClO <sub>3</sub> )
475.76	20	474.09	452.14	450.66	473.3	471.8	474.1	472.6	ν <sub>4</sub> , e δ <sub>as</sub> (ClO <sub>3</sub> )

[a] Most intensive matrix site, average of several measurements. [b] According to C<sub>3v</sub> symmetry. [c] Integrated relative infrared intensities, I(ν<sub>3</sub>) = 100.

experiments using different precursors and thermolysis temperatures, as well as by photolysis of the matrix using light of wavelength  $> 530$  nm, which causes a decrease in intensity for all  $\text{ClO}_3$  bands. The presence of  $\text{ClO}_3$  radicals with  $C_{3v}$  symmetry is evident from the following facts: i) the IR bands display an isotopic pattern (3:1), in accordance with a species, containing one chlorine atom; ii) in experiments with  $\text{Cl}_2^{18}\text{O}_4$  containing about 70%  $^{18}\text{O}$  four isotopomers are detected, consistent with a species containing three equivalent oxygen atoms; iii) four IR active fundamentals are observed for  $\text{ClO}_3$  ( $C_{3v}$ ) in accordance with the irreducible representation [Eq. (27)]:

$$\Gamma_{\text{vib}} = 2a_1 (\text{IR, Ra p}) + 2e (\text{IR, Ra dp}) \quad (27)$$

iv) for the mixed isotopomers  $\text{Cl}^{16}\text{O}_2^{18}\text{O}$  and  $\text{Cl}^{16}\text{O}^{18}\text{O}_2$  of  $C_s$  symmetry, the e modes are split into  $a'$  and  $a''$  components, as shown in Figure 5 (Table 5); and v) good agreement of observed IR data and predicted from quantum chemical calculations (see Table 6).

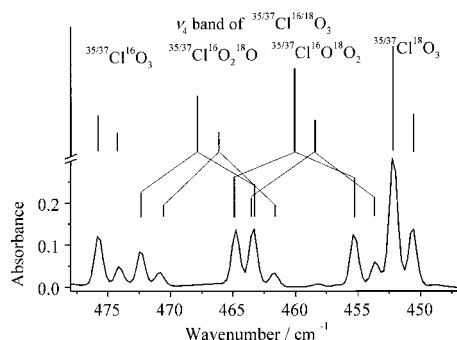


Figure 5. The  $\nu_4$  bands of an isotopic mixture of  $\text{ClO}_3$ . The components are compared with a simulated spectrum.

All these observations leave no doubt regarding the presence of  $\text{ClO}_3$  and confirm the assignment of the vibrational bands. In Figure 5 the  $\nu_4$  bands of the four isotopomers  $\text{Cl}^{16}\text{O}_3$ ,  $\text{Cl}^{16}\text{O}_2^{18}\text{O}$ ,  $\text{Cl}^{16}\text{O}^{18}\text{O}_2$  and  $\text{Cl}^{18}\text{O}_3$  were simulated using the molar ratio of 1:1.6:2.3:3.0, respectively. For the mixed species  $\text{Cl}^{16}\text{O}_2^{18}\text{O}$  and  $\text{Cl}^{16}\text{O}^{18}\text{O}_2$  the lower and higher wavenumber components are of  $a'$  symmetry, respectively, because the in plane atom in both cases is involved mostly in the  $a'$  mode. Also the overtones and combination modes of  $E$  symmetry show the same isotopic pattern as do the e-type fundamentals. Predicted IR spectra and geometric parameters of high level ab initio and DFT calculations are given in Table 6. The results agree quite well with the observed IR spectrum of  $\text{ClO}_3$  isolated in a Ne matrix, which resembles

Table 5. Vibrational wavenumbers [ $\text{cm}^{-1}$ ]<sup>[a]</sup> of mixed isotopomers of  $\text{ClO}_3$  isolated in Ne matrix.

$^{35}\text{Cl}^{16}\text{O}_2^{18}\text{O}$	$^{37}\text{Cl}^{16}\text{O}_2^{18}\text{O}$	$^{35}\text{Cl}^{16}\text{O}^{18}\text{O}_2$	$^{37}\text{Cl}^{16}\text{O}^{18}\text{O}_2$	assignment	$C_{3v}$	$C_s$
2148.4	2124.7	2113.9	2089.8	$2\nu_3$	E	$a''$
2124.7	2101.0	2131.8	2107.4			$a'$
1956.2	1941.9	1904.3	1889.2	$\nu_1 + \nu_3$	E	$a'$
1932.6	1917.6	1932.2	1917.6			$a'$
1635.9	1616.1	1597.2	1577.9	$\nu_2 + \nu_3$	E	$a''$
1616.1	1597.2	1620.6	1601.7			$a'$
1081.3	1069.4	1047.1	1034.9	$\nu_3$	e	$a''$
1060.6	1048.3	1071.3	1059.5			$a'$
888.3	885.5	873.5	870.7	$\nu_1$	$a_1$	$a'$
561.1	554.1	555.5	548.4	$\nu_2$	$a_1$	$a'$
472.4	470.9	464.8	463.4			$a''$
463.4	461.7	455.3	453.7	$\nu_4$	e	$a'$

[a] Most intensive matrix site.

closely the gas phase situation of the free radicals. The smallest deviations from experiment are seen for calculations, using UB3LYP and B3LYP methods.

**Evaluation of the molecular structure of  $\text{ClO}_3$ :** The experimental bond angle is calculated from the observed vibrational data by using the secular Equation [Eq. (28)]<sup>[193]</sup> for the symmetry class e in the point group  $C_{3v}$ .

$$|\mathbf{G}^* \mathbf{F}| = |\mathbf{E}^* \lambda| \quad (28)$$

For  $\mathbf{G}$  matrix: geometric parameters, masses;  $\mathbf{F}$  matrix: force constants;  $\mathbf{E}$ : unit matrix;  $\lambda$ : Eigenvalues. Under the assumption that the F matrix for each symmetry class is equal for the different isotopomers<sup>[193]</sup> and  $\lambda$  is proportional to  $\omega^2$  (harmonic wavenumber) the following Equation (29) is valid.

$$\frac{\prod_i \omega_i^2}{\prod_i \omega_i'^2} = \left| \frac{\mathbf{G}}{\mathbf{G}'} \right| \quad (29)$$

This Equation is a special form of the Redlich–Teller product rule.<sup>[193]</sup> By using the matrix elements  $G_{ij}$ <sup>[194]</sup> and calculating for different angles the six possible quotients of the determinants  $\mathbf{G}$  and  $\mathbf{G}'$  for the four isotopomers  $^{35}\text{Cl}^{16}\text{O}_3$ ,  $^{37}\text{Cl}^{16}\text{O}_3$ ,  $^{35}\text{Cl}^{18}\text{O}_3$ , and  $^{37}\text{Cl}^{18}\text{O}_3$  it is shown, that only the quotients in the symmetry class e are dependent of the OClO bond angle. Figure 6 shows two of the six plots  $\det|\mathbf{G}/\mathbf{G}'|$  against the OClO bond angle (two with a positive slope and four with a negative slope) in the symmetry class e. The values for the quotients of the square products of the observed wavenumbers obtained from the isotopic pairs  $^{35}\text{Cl}^{16}\text{O}_3/^{37}\text{Cl}^{16}\text{O}_3$  and  $^{35}\text{Cl}^{16}\text{O}_3/^{37}\text{Cl}^{18}\text{O}_3$  lead to the bond angles 110.3 and 130.6°, respectively. The large discrepancy between both results is

Table 6. Calculated vibrational wavenumbers [ $\text{cm}^{-1}$ ], band intensities ( $\text{km mol}^{-1}$ , in italic), bond length [pm] and bond angle [°] of  $^{35}\text{Cl}^{16}\text{O}_3$ .

Method	MP2 <sup>[a]</sup>	RMP2 <sup>[b]</sup>	B3LYP <sup>[c]</sup>	UB3LYP <sup>[d]</sup>	B3LYP <sup>[e]</sup>	SVWN <sup>[f]</sup>	expt <sup>[g]</sup>
mode e		1180 (77)	1050 (110)	1081	1080	1095	1081 (100)
$a_1$		987 (2)	903 (4)	924	923	921	905 (3)
$a_1$		581 (21)	551 (21)	564	564	548	567 (11)
e		479 (16)	464 (12)	472	471	450	476 (20)
$r$ (Cl–O)	148.1	149.2	149.2	145.2	145.4	146.1	148.5 ± 2
$\angle$ (OClO)	114.6	113.8	113.8	114.2	114.1	114.1	113.8 ± 1

[a] Ref. [116]. [b] Ref. [117]. [c] /cc-pVQ(5)Z, Ref. [40]. [d] Ref. [39]. [e] Ref. [118]. [f] Ref. [38]. [g] This work, relative integrated band intensities in italic.



due to the anharmonic nature of the experimental wavenumbers. They can be corrected by introduction of the anharmonic constants  $\delta_i$  according to Equation (30).

$$\frac{\prod_i \omega_i^2}{\prod_i \omega_i'^2} = \left( \frac{\nu_3 \cdot (1 + \delta_3)}{\nu_3' \cdot (1 + \delta_3')} \right)^2 \cdot \left( \frac{\nu_4 \cdot (1 + \delta_4)}{\nu_4' \cdot (1 + \delta_4')} \right)^2 = \frac{\prod_i \nu_{\text{corr.}}^2}{\prod_i \nu_{\text{corr.}}'^2} = \left| \frac{G}{G'} \right| \quad (30)$$

With the approximation:  $\nu_{\text{obs}} \cdot (1 + \delta) - \nu_{\text{obs}}' \cdot (1 + \delta') = \Delta \nu_{\text{obs}} \cdot (1 + \delta'') = \Delta \omega$ ,<sup>[195]</sup> Equation (30) can be rearranged to Equation (31).

$$\frac{\prod_i \omega_i^2}{\prod_i \omega_i'^2} = \left( \frac{\nu_3}{\nu_3 - \Delta \nu_3 \cdot (1 + \delta_3'')} \right)^2 \cdot \left( \frac{\nu_4}{\nu_4 - \Delta \nu_4 \cdot (1 + \delta_4'')} \right)^2 = \frac{\prod_i \nu_{\text{corr.}}^2}{\prod_i \nu_{\text{corr.}}'^2} = \left| \frac{G}{G'} \right| \quad (31)$$

As seen in Figure 6 each ratio of determinants can be described as an approximate linear function of the bond angle  $\alpha$  [Eq. (32)].

$$\left| \frac{G}{G'} \right| = m \cdot \alpha + b \leftrightarrow \left( \frac{\prod_i \nu_{\text{corr.}}^2}{\prod_i \nu_{\text{corr.}}'^2} - b \right) \cdot m^{-1} = \alpha \quad (32)$$

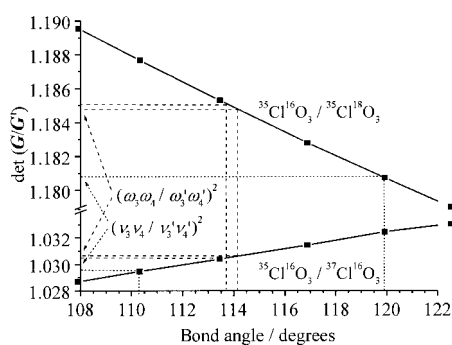


Figure 6. Plots of  $|G/G'|$  matrix element ratios in the representation  $e$  for different  $\text{ClO}_3$  isotopomers versus the OCIO bond angle. Pointed lines indicate the observed  $(\nu_3 \cdot \nu_4 / \nu_3' \cdot \nu_4')^2$  values. After a proper anharmonicity correction (see text) the dashed lines lead to the  $\alpha_e$  bond angle (113.8°) of  $\text{ClO}_3$ .

By varying  $\delta_3''$  and  $\delta_4''$ , in principle a corrected bond angle can be obtained, although there are several possible solutions. Taking into account that the potential of a deformation mode in general is less anharmonic than a stretching mode (due to the non-dissociative character of the deformation mode) one can assume a  $\delta_4'' / \delta_3''$  ratio in the range of 0 to 0.25. With this assumption matching of the bond angle for two  $|G/G'|$  slopes can be obtained by varying  $\delta_3''$  (dashed lines in Figure 6). By using all eight combinations of the six quotients (2 positive slope  $\times$  4 negative slope) an average bond angle of  $113.8 \pm 0.7^\circ$  is obtained. This bond angle should resemble the  $\alpha_e$  value for  $\text{ClO}_3$  in the gas phase within  $\pm 1^\circ$ , because it has been shown for the similar OCIO radical that the bond angle determined from Ne matrix vibrational data<sup>[5]</sup> is identical within the estimated error limit of  $\pm 0.2^\circ$  to the  $\alpha_e$  gas phase value.<sup>[196]</sup>

With the bond angle of  $\text{ClO}_3$  obtained as discussed above, calculations of the general valence force field using the NCA program<sup>[197]</sup> are performed. As input data all wavenumbers of the four fundamentals of eight isotopomers (Tables 4 and 5) are used after the required anharmonic corrections are performed according to Becher.<sup>[195]</sup> All data are iterated simultaneously until minimum least square deviation between observed and calculated vibrational data are obtained. The result is presented in Table 7. Due to the large number of

Table 7. Force constants ( $10^2 \text{ N m}^{-1}$ , normalized on 100 pm bond length) and potential energy distribution (PED) of  $\text{ClO}_3$ .

	Force constants		PED		
	NCA	$\nu_3$	$\nu_1$	$\nu_2$	$\nu_4$
$f_r$	6.780	1.00	0.96		
$f_{rr}$	0.156				
$f'_{ra}$	0.050				
$f_{ra}$	0.228				
$f_a$	2.472			0.52	1.74
$f_{aa}$	1.103			0.46	-0.78

independent vibrational isotopic input data ( $>10$ ) all six inner-force constants are overestimated and should be reliable. This is a rare case and comparison with ab initio values would be desirable. However, to our knowledge such data are not available from the literature at present. Most striking is the unusual large interaction force constant  $f_{aa}$  which causes a large potential energy distribution (PED) for  $\nu_2$  and  $\nu_4$ . In other molecules  $f_{aa}$  may be much smaller. Hence in the  $\text{ClO}_3$  radical the OCIO bending causes a strong electron flow to the neighbouring bonds. Use of the stretching force constant  $f_r$  offers an estimate of the ClO bond length. For this purpose ClO force constants of different species containing ClO bonds (Table 8) are plotted against their respective bond length and fitted with an exponential function (Figure 7).

Table 8. Structural parameters (pm, °,  $10^2 \text{ N m}^{-1}$ ) for different species containing ClO bonds.

	$\text{Cl}_2\text{O}^{[a]}$	$\text{HOCl}^{[b]}$	$\text{ClO}^{[c]}$	$\text{FCIO}^{[d]}$	$\text{OCIO}^{[e]}$	$\text{ClClO}_2^{[f]}$	$\text{FCIO}_2^{[g]}$	$\text{FCIO}_3^{[h]}$
$r_e$ (ClO)	169.59	168.91	156.96	148.4	146.98	143.68	141.99	139.7
$\angle$ OCIO	–	–	–	–	117.4	114.63	115.03	115.8
$f$ (ClO)	2.959	3.79	4.71	7.04	7.03	8.52	9.36	10.13

[a] Ref. [198–200]. [b] Ref. [201, 202]. [c] Ref. [81, 203]. [d] Ref. [203]. [e] Ref. [19]. [f] Ref. [204]. [g] Ref. [205, 206]. [h] Ref. [203, 207].

Assuming a 90% confidence level, the ClO bond length is estimated to be  $148.5 \pm 2$  pm. The estimated structural parameters of  $\text{ClO}_3$  from vibrational data, termed “experimentally” are compared in Table 6 with calculated parameters. Our “experimental” structure is within the error limits in accordance with most theoretical predictions. A future accurate structural determination either by microwave or/and high resolution IR spectroscopy is needed, to decide, which prediction is optimal.

**The UV/Vis spectrum of  $\text{ClO}_3$ :** In order to record UV spectra of  $\text{ClO}_3$  with less interference from by-products as reported in

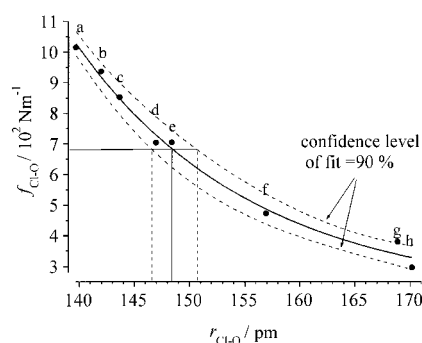


Figure 7. Correlation of ClO force constants versus ClO bond lengths for different ClO species. With the ClO force constants of ClO<sub>3</sub> 6.78 (see text) a ClO bond length of 148.5 pm is estimated. Under consideration of a 90% confidence level of the exponential fit the error is estimated to be  $\pm 2$  pm. a) FClO<sub>3</sub>, b) FClO<sub>2</sub>, c) ClClO<sub>2</sub>, d) ClO<sub>2</sub>, e) FClO, f) ClO, g) ClOH, h) ClOCl.

our previous note,<sup>[2]</sup> we mainly use FOCIO<sub>3</sub> as source for ClO<sub>3</sub> in this study. After low pressure flash thermolysis of FOCIO<sub>3</sub> with subsequent isolation of the products in a Ne matrix, the IR spectrum shows FOCIO<sub>3</sub> and ClO<sub>3</sub> along with the by-products FO, FOO, FOOF and ClO<sub>2</sub>. Hence it can be assumed that the reaction proceeds as follows [Eqs. (33)–(37)].

primary:



secondary:



A crude UV/Vis spectrum averaged from the five “best” flash thermolysis experiments has been recorded, as well as the reference UV spectra of OCIO and FOCIO<sub>3</sub>. The unstructured spectra of FOO, and FOOF, taken from literature,<sup>[208, 209]</sup> are added to the reference spectrum of FOCIO<sub>3</sub> according to the estimated portions in the thermolysis experiments. This leads to a synthetic reference spectrum in the 200–270 nm range. Applying both reference spectra for computational spectra subtraction of the by-products, to the raw spectrum, the spectrum presented in Figure 8 is obtained. Smoothing of the absorption profile and fitting both bands with Gauss functions leads to the spectrum depicted in Figure 9.

The high energy band is found at  $\nu_{\text{max}} = 32100 \text{ cm}^{-1}$  (311.5 nm), a half width of  $6800 \text{ cm}^{-1}$ , an onset (0–0 transition) at  $\approx 20000 \text{ cm}^{-1}$ , and the low energy band at  $\nu_{\text{max}} = 23150 \text{ cm}^{-1}$  (432.0 nm), half width  $6400 \text{ cm}^{-1}$ , onset at  $\approx 13300 \text{ cm}^{-1}$ . The band intensity ratio is 1:0.7. These results are in good agreement with a recent photoelectron study of the ClO<sub>3</sub><sup>-</sup> anion<sup>[25]</sup> and the theoretically predicted electronic transitions of ClO<sub>3</sub>.<sup>[210]</sup> In the photoelectron spectroscopy study of ClO<sub>3</sub><sup>-</sup> with a 157 nm laser excitation three electronic states ( $X^2A_1$ ,  $A^2A_1$ ,  $B^2E$ ) have been probed. The above-

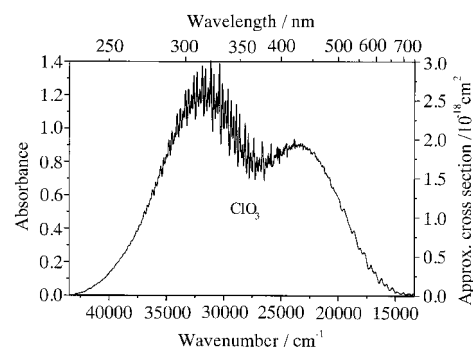


Figure 8. UV/Vis spectrum of ClO<sub>3</sub> isolated in a neon matrix. The radical is produced by low pressure thermolysis of FOCIO<sub>3</sub>. Bands of by-products are subtracted by using reference spectra (see text).

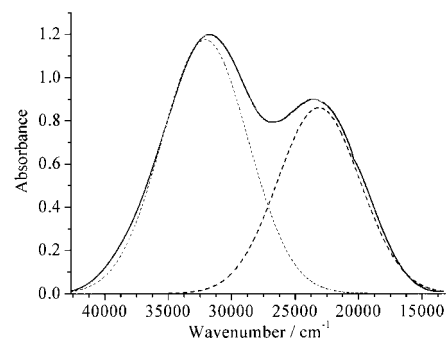


Figure 9. Smoothed UV/Vis spectrum of ClO<sub>3</sub> presented in Figure 8. Dashed lines are simulated Gaussian profiles of the bands.

mentioned maxima of the electronic transition in the UV/Vis spectrum of ClO<sub>3</sub> can be correlated with the energy gaps between the adiabatic value of the ground state ( $X^2A_1$ ) and the vertical detachment energies of the excited states. For the  $A^2A_1 \leftarrow X^2A_1$  and  $B^2E \leftarrow X^2A_1$  transition, this gap amounts to  $1.5 \pm 0.2 \text{ eV}$  ( $12100 \text{ cm}^{-1}$ ) and  $2.85 \pm 0.2 \text{ eV}$  ( $23000 \text{ cm}^{-1}$ ), respectively. The first excited state has a transition moment equal to zero<sup>[210]</sup> and is not observed experimentally. The second transition is in excellent agreement with the observed UV absorption. In a recent high level ab initio study of ClO<sub>3</sub>,<sup>[210]</sup> using several different methods the vertical excitation energies into the  $A^2A_2$ ,  $B^2E$ , and  $C^2E$  states are predicted to be in the range  $1.25\text{--}1.49 \text{ eV}$  ( $\approx 11000 \text{ cm}^{-1}$ ),  $2.55\text{--}2.90 \text{ eV}$  ( $\approx 22000 \text{ cm}^{-1}$ ) and  $3.55\text{--}3.93 \text{ eV}$  ( $\approx 30200 \text{ cm}^{-1}$ ) with a relative transition moment for the last two transitions of about 1.2:1, respectively. Both calculated band maxima are on average a little too low and the band intensities are reversed, in comparison to the observed UV spectrum.

As demonstrated in the preceding chapters the UV spectra of ClO and OCIO, isolated in Ne matrix, are in good agreement with their respective gas phase spectra. They are blue-shifted by only a few tenth of a nm; some minor effects caused by the rotational fine structure and hot bands in the gas-phase spectra are quenched in the matrix. Hence it can be assumed that the UV spectrum of ClO<sub>3</sub> depicted in Figure 8 corresponds very well to the yet unknown gas phase spectrum. The cross section scale placed on the right axis in Figure 8 is deduced from photolysis experiments. Photolysis with light  $\lambda > 530 \text{ nm}$  causes a decrease of the ClO<sub>3</sub> absorption and a

simultaneous increase of ClO absorption. With help of the generally accepted gas phase cross section of ClO ( $\sigma_{\max} = 5.2 \times 10^{-18} \text{ cm}^2$  at  $\lambda_{\max} = 260 \text{ nm}$ <sup>[211]</sup>) the cross section of ClO<sub>3</sub>  $\sigma = 3.0 (\pm 0.5) \times 10^{-18} \text{ cm}^2$  at the strongest progression band (11,0,0) at  $\lambda = 321.15 \text{ nm}$  are adjusted.

The high-energy UV band, which is shown in Figure 10 on an expanded scale, exhibits a pronounced vibrational fine structure. The most intense peaks in this fine structure form the main progression and are separated by about  $800 \text{ cm}^{-1}$ . Each peak in this progression is accompanied by two satellites

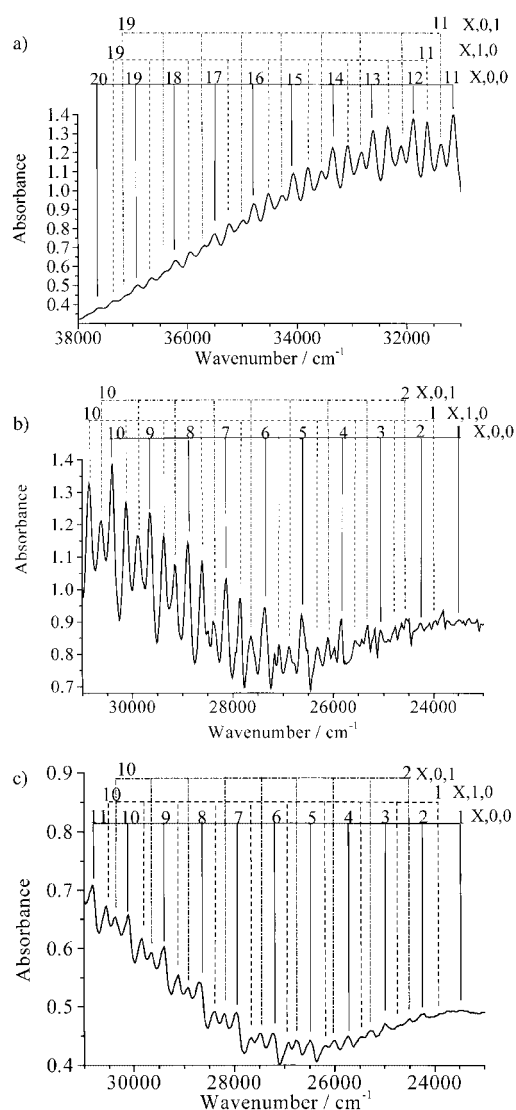


Figure 10. a) and b) Structured part of the UV/Vis spectrum of ClO<sub>3</sub> in expanded scale from Figure 8. c) Part b) for Cl<sup>18</sup>O<sub>3</sub>.

due to two additional vibrations in the excited C<sup>2</sup>E state. Due to the antibonding nature of this excited state<sup>[210]</sup> all fundamentals should be lower than in the ground state (1081, 905, 567, 476  $\text{cm}^{-1}$ ). According to the selection rules, excitation of all fundamentals in the C<sup>2</sup>E ← X<sup>2</sup>A<sub>1</sub> transition is allowed.<sup>[212]</sup> A shifted asymmetric stretching mode (from  $\approx 1080$  to  $800 \text{ cm}^{-1}$ ) seems to be the origin of the main progression, because its normal coordinate correlates well with the dissociation path into OClO + O,<sup>[210]</sup> and the relative isotopic

vibrational shifts of  $\Delta\nu(\text{Cl}^{16/18}\text{O}_3)/\nu(\text{Cl}^{16}\text{O}_3)$  which are equal in the ground and excited state. The first high energy satellite on each progression bands could be either a  $\approx 250$  or a  $\approx 1050 \text{ cm}^{-1}$  vibration depending whether it “sits” on the quantum number  $n$  or  $(n - 1)$  of the progression, respectively. Due to the choice of the asymmetric stretching vibration in the main progression and the antibonding nature of the excited state, we assume the first satellite “sits” on the quantum number  $n$  (blue-shifted by  $250 \text{ cm}^{-1}$ ) and can be correlated with the symmetric deformation mode in the ground state ( $567 \text{ cm}^{-1}$ ). The second high energy satellite (blue-shifted by  $500 \text{ cm}^{-1}$ ) can be correlated with the  $905 \text{ cm}^{-1}$  symmetrical stretching mode in the ground state.

As seen in Figure 10b, c the origin of the progression somewhere below  $24000 \text{ cm}^{-1}$  is uncertain. To determine the origin of the high energy UV band, the spectrum of Cl<sup>18</sup>O<sub>3</sub> is measured (Figure 10c). The isotopic shifts between the bands of Cl<sup>16</sup>O<sub>3</sub> and Cl<sup>18</sup>O<sub>3</sub> are proportional to the quantum numbers and should be nearly zero at the origin of the transition. By fitting the linear difference of equivalent bands of Cl<sup>16</sup>O<sub>3</sub> and Cl<sup>18</sup>O<sub>3</sub> against arbitrary quantum numbers, the (0,0,0) transition is finally located at about  $22700 \text{ cm}^{-1}$ . With the true quantum numbers available, the whole progression spectra of both isotopomers are simulated with the second order approximation of Herzberg<sup>[213]</sup> [Eq. (38)].

$$\nu_{\text{obs}} = \nu_{00} + \sum_i \omega_i^0 \nu_i + \sum_i \sum_{k \geq 1} \chi_{ik}^0 \nu_i \nu_k; \omega_i^0 = \omega_i + \chi_{ii} + 0.5 \sum_{i+k} \chi_{ik}^0; \chi_{ik}^0 = \chi_{ik} \quad (38)$$

with  $\omega_i^0$  = harmonic wavenumber and  $\chi_{ik}^0$  = anharmonic constants. The observed wavenumbers are taken directly from the band maxima or they are examined from their second derivative spectra. For the simulation of the Cl<sup>16</sup>O<sub>3</sub> data, the three harmonic wavenumbers and five anharmonicity constants were iterated simultaneously until the sum of deviations reached a minimum. Subsequently the anharmonicity constants are kept nearly constant for the simulation of the Cl<sup>18</sup>O<sub>3</sub> data. The results of the simulation are presented in Tables 9 and 10. In general the deviation between observed and calculated band position in Table 9 are in the order of the uncertainty limits of the measured band positions. Hence the resulting parameters presented in Table 10 appear to be reliable. The anticipated much lower vibrations in the excited C<sup>2</sup>E state and their assignments in comparison to the ground state are supported by their equivalent isotopic effects.

There is additional structure in the spectrum which may be due to matrix site splitting, a further vibration or the splitting of the e modes. Definitive conclusions, however, are not possible, because residual incompenations as a result of the spectra subtraction which causes some artefacts. The analysis of the fine structure of the low energy band does not lead to reasonable results, due to the overlapping of bands. It can be assumed that there is a main progression with several satellites with a band separation by  $900$  to  $1000 \text{ cm}^{-1}$ .

A critical evaluation of the literature data for other claims for the UV/Vis spectrum of ClO<sub>3</sub><sup>[96, 97, 120, 121, 140]</sup> leads to the conclusion that in all cases something else was observed, but not ClO<sub>3</sub>. In a recent study on laser photolysis of peroxydisulfate solutions containing chlorate, a very broad unsym-

Table 9. Assignment for the progressions of Ne matrix isolated ClO<sub>3</sub> in the C<sup>2</sup>E ← X<sup>2</sup>A<sub>1</sub> transition.

Assignment			Cl <sup>16</sup> O <sub>3</sub>		Cl <sup>18</sup> O <sub>3</sub>		assignment			Cl <sup>16</sup> O <sub>3</sub>		Cl <sup>18</sup> O <sub>3</sub>	
ν <sub>3</sub>	ν <sub>1</sub>	ν <sub>2</sub>	observed	diff. <sup>[a]</sup>	observed	diff. <sup>[a]</sup>	ν <sub>1</sub>	ν <sub>2</sub>	ν <sub>3</sub>	observed	diff. <sup>[a]</sup>	observed	diff. <sup>[a]</sup>
				x [cm <sup>-1</sup> ]							x [cm <sup>-1</sup> ]		
0	0	0	22696		22696								
1	0	0	23482	-6	23441	-18	12	0	0	31883	2	31536	-6
1	0	1			23724	0	12	0	1	32098	-8	31776	10
1	1	0	23964	-16			12	1	0	32347	0	31974	-6
2	0	0	24269	-6	24207	-11	13	0	0	32616	1	32253	5
2	0	1	24545	2	24486	7	13	0	1	32825	-11	32483	14
2	1	0	24764	-1	24663	-17	13	1	0	33085	6	32680	-3
3	0	0	25063	6	24970	-2	14	0	0	33350	5		
3	0	1	25320	-1	25233	3	14	0	10	33551	-10		
3	1	0	25562	18	25430	-1	14	1	0	33807	1		
4	0	0	25840	5	25720	-1	15	0	0	34066	-3		
4	0	1	26103	9	25991	16	15	0	1	34275	-7		
4	1	0	26319	-1	26164	-14	15	1	0	34524	-4		
5	0	0	26635	28	26465	-1	16	0	0	34995	-2		
5	0	1	26881	19	26720	4	16	0	1	35236	-9		
5	1	0	27089	-1	26914	-6	16	1	0	35236	-9		
6	0	0	27367	-8	27205	0	17	0	0	35499	-5		
6	0	1	27640	14	27463	11	17	0	1	35707	-1		
6	1	0	27847	-8	27657	0	17	1	0	35958	0		
7	0	0	28129	-9	27946	6	18	0	0	36231	17		
7	0	1	28393	8	28188	5	18	0	10	36410	-4		
7	1	0	28612	-4	28381	-9	18	1	0	36670	4		
8	0	0	28898	2	28674	4	19	0	0	36910	-10		
8	0	1	29138	-1	28910	1	19	0	1	37120	5		
8	1	0	29369	-3	29112	-	19	1	0	37369	0		
9	0	0	29656	6	29394	-1	20	0	0	37643	23		
9	0	1	29881	-7	29643	12	20	0	1	37814	3		
9	1	0	30125	3	29841	1	20	1	0	38073	6		
10	0	0	30400	2	30111	-5	21	0	0	38314	-2		
10	0	1	30623	-9	30358	11	21	1	1	38513	-11		
10	1	0	31864	-5	30553	-5	21	1	0	38760	0		
11	0	0	31138	-4	30821	-10	22	0	0	38986	-21		
11	0	1	31373	2	31065	6				Σ x  = 235 <sup>[b]</sup>		Σ x  = 139 <sup>[c]</sup>	
11	1	0	31626	16	31269	-2							

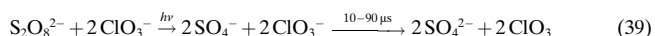
[a] Observed – calculated = difference. [b] Quantum number 6 to 19 included. [c] Quantum number 6 to 13 included.

Table 10. Wavenumbers [cm<sup>-1</sup>] and anharmonic constants of the normal modes of Cl<sup>16</sup>O<sub>3</sub>/Cl<sup>18</sup>O<sub>3</sub> in the excited C<sup>2</sup>E state.

Wavenumber <sup>[a]</sup>	Assign- ment	Anharmonic constant			
794 (1081)	ω <sub>1</sub> <sup>0</sup> e ν <sub>as</sub>	χ <sub>11</sub>	-2.4		
766 (1047)					
498 (905)	ω <sub>2</sub> <sup>0</sup> a <sub>1</sub> ν <sub>s</sub>	χ <sub>12</sub>	-2.4	χ <sub>22</sub>	-3.7
470 (860)					
280 (567)	ω <sub>3</sub> <sup>0</sup> a <sub>1</sub> δ <sub>s</sub>	χ <sub>13</sub>	≈ -4.3	χ <sub>23</sub>	0
272 (550)			≈ -3.7		χ <sub>33</sub> -3.4

[a] In parenthesis ground state values, second values belong to Cl<sup>18</sup>O<sub>3</sub>.

metrical band with a maximum at about 330 nm and a flat band at about 450 nm has been reported.<sup>[122]</sup> The shape of this spectrum resembles in principle that of the ClO<sub>3</sub> spectrum in Figure 8, so that the following reaction sequence may have occurred [Eq. (39)].



In a similar way, the UV/Vis spectrum of the isoelectronic SO<sub>3</sub><sup>-</sup> radical anion has been measured.<sup>[214]</sup> In comparison to the UV/Vis spectrum of ClO<sub>3</sub> the spectrum is structureless and blue-shifted by about 50 nm with a cross section of 4.0 ×

10<sup>-18</sup> cm<sup>2</sup> at 260 nm and a low energy band at 360 nm of much lower intensity, than the respective band of ClO<sub>3</sub>.

**The photochemistry of ClO<sub>3</sub>:** Photolysis of ClO<sub>3</sub> isolated in a Ne matrix with a tungsten halogen lamp in combination with a cut-off filter λ > 455 nm leads to ClO (<sup>35</sup>ClO 844.0 and 839.5 cm<sup>-1</sup>) and O<sub>2</sub> (1549.1 cm<sup>-1</sup>) while ClO<sub>3</sub>, isolated in a O<sub>2</sub> matrix, leads mainly to O<sub>3</sub> (1035.5, 702.4 cm<sup>-1</sup>) and OCIO (1104.7, 942.1, 453.4 cm<sup>-1</sup>) after the same irradiation. In the former case the bands of ClO and O<sub>2</sub> are slightly shifted in comparison to the direct matrix isolated species (<sup>35</sup>ClO 843.8; O<sub>2</sub> 1554 cm<sup>-1</sup>). The primary step after excitation into the B<sup>2</sup>E state is dissociation [Eq. (40)].



In Ne matrix, reaction of the oxygen atom with OCIO occurs in the matrix cage according to Equation (41).



In O<sub>2</sub> matrix the O atom reacts with one of the surrounding O<sub>2</sub> molecules and O<sub>3</sub> is formed. Two different ClO·O<sub>2</sub> complexes are formed in a Ne matrix as indicated by the two IR bands at 844.0 and 839.5 cm<sup>-1</sup>. In addition a new UV

absorption on the short wavelength wing of the ClO band appears (Figure 11).

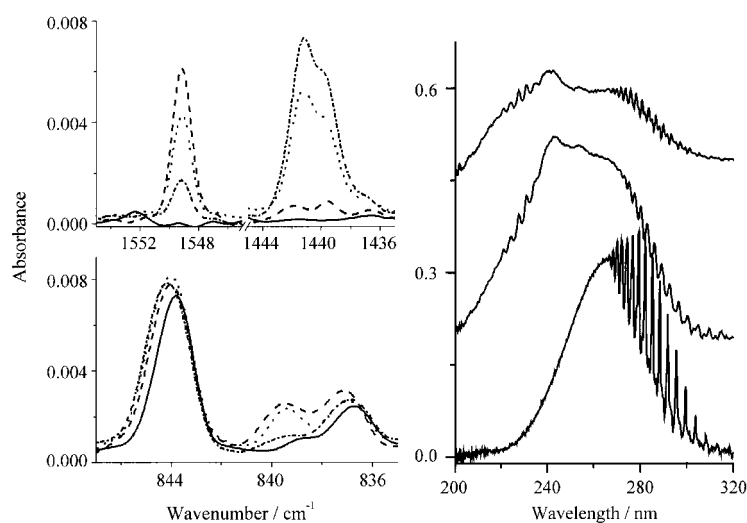


Figure 11. IR and UV spectra of the low pressure thermolysis products of ClOClO<sub>3</sub> isolated in a neon matrix. (—): IR spectrum of monomeric <sup>35</sup>ClO (843.8 cm<sup>-1</sup>) <sup>37</sup>ClO (836.8 cm<sup>-1</sup>) in the <sup>2</sup>Π<sub>3/2</sub> state and <sup>35</sup>ClO (838.5 cm<sup>-1</sup>) in the <sup>2</sup>Π<sub>1/2</sub> state. Lower trace; right: UV spectrum of ClO after subtraction of the bands of all by-products (compare with Figure 1). (---): IR-spectrum after visible light photolysis (λ > 455 nm) of about 80% of the matrix isolated ClO<sub>3</sub> yielding <sup>35</sup>ClO·O<sub>2</sub> (844.0 cm<sup>-1</sup>) <sup>37</sup>ClO·O<sub>2</sub> (837.0 cm<sup>-1</sup>) in the <sup>2</sup>Π<sub>3/2</sub> state and <sup>35</sup>ClO·O<sub>2</sub> (839.5 cm<sup>-1</sup>) in the <sup>2</sup>Π<sub>1/2</sub> state. At 1549.1 cm<sup>-1</sup> the O<sub>2</sub> stretching band of ClO·O<sub>2</sub> in the <sup>2</sup>Π<sub>1/2</sub> state is shown. Middle trace, right: difference UV spectrum (photolysis minus thermolysis) of ClO·O<sub>2</sub>. In contrast to ClO obtained by thermolysis (lower trace) the higher populated <sup>2</sup>Π<sub>1/2</sub> state is clearly shown. . . . and ●—: IR spectra after 255 nm photolysis for 10 and 30 min. Mainly the complex ClO·O<sub>2</sub> (<sup>2</sup>Π<sub>1/2</sub>) disappeared (839.5 cm<sup>-1</sup>) and at 1440 cm<sup>-1</sup> ClOO appeared. Upper trace: difference UV spectrum (photolysis 455 nm minus photolysis 255 nm) indicating mainly the loss of ClO·O<sub>2</sub> (<sup>2</sup>Π<sub>1/2</sub>).

Further photolysis of the Ne matrix at 255 nm (high-pressure Hg lamp + cut-off filter) causes a decrease of the IR bands at 839.5 and 1549.1 cm<sup>-1</sup>, a small increase of the band at 844.0 cm<sup>-1</sup>, appearance of a new band at 1441 cm<sup>-1</sup> (ClOO) as well as disappearance of the new UV band. The change of the UV spectrum is displayed in Figure 11. These findings can be interpreted as follows: On photolysis of ClO<sub>3</sub> in Ne matrix, a complex of the composition ClO·O<sub>2</sub> in analogy to the FO·O<sub>2</sub> complex<sup>[215]</sup> is formed, which is characterized by two UV bands (see Figure 11) and an IR band at 839.5 cm<sup>-1</sup>. Additionally, a second ClO–O<sub>2</sub> species is formed, which is assigned to a ClO radical situated next to an oxygen molecule in the same Ne matrix cage showing a weaker interaction. The former complex (839.5 cm<sup>-1</sup>) is more photosensitive at 255 nm, than the latter ClO species (844.0 cm<sup>-1</sup>) and is depleted into several channels: i) It is converted into the weak ClO·O<sub>2</sub> complex (844.0 cm<sup>-1</sup>, maybe the isomer OCl·O<sub>2</sub>) with an IR inactive O<sub>2</sub> molecule; ii) an oxygen atom left the matrix cage and some ClOO (1441 cm<sup>-1</sup>) is formed and iii) it dissociates into O<sub>2</sub>+Cl+O. In the difference UV spectrum before and after photolysis at 255 nm, mainly a peroxy radical band at ≈260 nm with a vibrational progression of ≈670 cm<sup>-1</sup> and, in the ClO region, the <sup>2</sup>Π<sub>1/2</sub> subsystem are clearly seen. Hence, we interpret the new UV and IR (839.5, 1549.1 cm<sup>-1</sup>) absorber as a ClO·O<sub>2</sub> complex with ClO in the excited <sup>2</sup>Π<sub>1/2</sub> state.

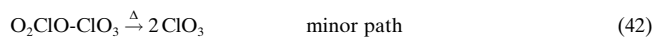
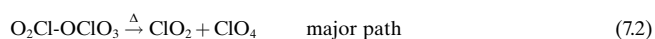
There is another interesting observation: Visible light photolysis of Ne matrix isolated ClO<sub>3</sub> causes an intensity increase of the weak bands of Cl<sub>2</sub>O<sub>6</sub>, which are formed during the deposition process. Presumably a small fraction of (ClO<sub>3</sub>)<sub>2</sub> van der Waals dimers are activated in this manner and asymmetric bond linking of (ClO<sub>3</sub>)<sub>2</sub> to O<sub>2</sub>ClOClO<sub>3</sub> occurs.

The experimentally observed threshold for slow (5% in one hour) photodissociation of ClO<sub>3</sub> is at 610 nm. With more intense light sources, the threshold may be observable at the onset of the absorption band at ≈750 nm. Interestingly the energy of the onset (≈160 kJ mol<sup>-1</sup>) is within the uncertainty, identical to the first Cl–O bond energy of ClO<sub>3</sub> calculated by Francisco<sup>[118]</sup> and Janoschek.<sup>[40]</sup>

## Chlorine Tetroxide

In our preliminary study of matrix isolated ClO<sub>4</sub> radicals IR and UV/Vis spectra are recorded and analysed in some detail.<sup>[3]</sup> Based on the five observed IR bands C<sub>3v</sub> symmetry for the radical was assumed. However, several recent theoretical studies disagree with this conclusion and C<sub>2v</sub> symmetry is predicted.<sup>[38–40, 133, 134]</sup> We want to report here on an extended spectroscopic study of this elusive radical, in order to establish an experimental basis for final conclusions.

**The IR spectrum of ClO<sub>4</sub>:** Dichlorine hexoxide, Cl<sub>2</sub>O<sub>6</sub>, is the best precursor for the synthesis of ClO<sub>4</sub> radicals and is used throughout in this study. Low pressure flash pyrolysis of Cl<sub>2</sub>O<sub>6</sub>, highly diluted in a matrix gas in front of the matrix support with subsequent quenching of the products at low temperatures, leads to ClO<sub>2</sub>, ClO<sub>4</sub>, ClO<sub>3</sub> and minor amounts of HClO<sub>4</sub>. The formation of these products is assumed to proceed as follows [Eqs. (7.2, 42–45)].



secondary:



Neon matrix reference spectra of all by-products and Cl<sub>2</sub>O<sub>6</sub> are recorded for use of computational spectra subtraction. After “treatment” of the ten best raw neon matrix spectra of the Cl<sub>2</sub>O<sub>6</sub> pyrolysis products, with these reference spectra the resulting spectra are averaged and smoothed. The final spectrum is presented in Figure 12.

It shows three further weak IR absorptions in addition to the five known bands of ClO<sub>4</sub>.<sup>[3]</sup> Since the relative intensities of the new bands are the same in all experiments and since they decrease uniformly after photolysis of the matrix (tungsten-halogen lamp, cut off filter λ > 495 nm), all bands should belong to ClO<sub>4</sub>. Further support for the identity of the new bands and their assignment come from experiments with isotopically enriched Cl<sub>2</sub><sup>18</sup>O<sub>6</sub>. In the spectral regions of the

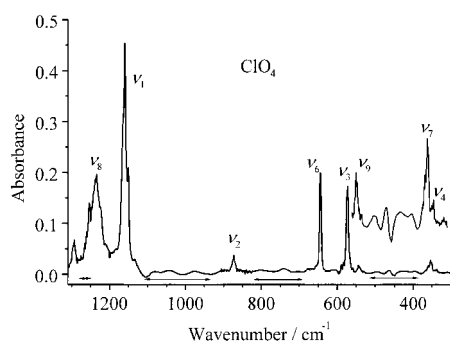


Figure 12. IR spectrum of ClO<sub>4</sub> isolated in a neon matrix. Bands of residual Cl<sub>2</sub>O<sub>6</sub> and of by-products (Cl<sub>2</sub>O<sub>4</sub>, ClO<sub>2</sub>, HClO<sub>4</sub>) are subtracted and the regions are indicated by arrows.

fundamentals the eight observed IR bands, including their Cl<sup>18</sup>O<sub>4</sub> counterparts, are presented in Figure 13 in expanded scale. All vibrational band positions and relative band intensities are gathered in Table 11.

With the increased number of IR active fundamentals, the ClO<sub>4</sub> radical must exhibit a lower symmetry than C<sub>3v</sub> as previously assumed.<sup>[3]</sup> Depending on the possible symmetries C<sub>2v</sub> or C<sub>s</sub> the nine vibrational motions of ClO<sub>4</sub> are divided into the irreducible representations [Eqs. (45) and (46)].

$$C_{2v}: \Gamma_{\text{vib}} = 4a_1 (\text{IR, Ra p}) + a_2 (-) + 2b_1 (\text{IR, Ra dp}) + 2b_2 (\text{IR, Ra dp}) \quad (45)$$

with 8 IR active fundamentals or

$$C_s: \Gamma_{\text{vib}} = 6a' (\text{IR, Ra p}) + 3a'' (\text{IR, Ra dp}) \quad (46)$$

with nine IR active fundamentals. With the experimental data alone, one can not distinguish unambiguously between C<sub>2v</sub> and C<sub>s</sub> symmetry, because the activated a<sub>2</sub> mode in C<sub>s</sub> symmetry can be very weak. To achieve a decision between C<sub>2v</sub> and C<sub>s</sub>, the vibrational data in Table 11 are compared with the theoretical predicted data in Tables 12 and 13. According

to most theoretical calculations, the ClO<sub>4</sub> radical should possess C<sub>2v</sub> symmetry, with two short and two long ClO bonds forming bond angles of about 114 and 95°, respectively (see Table 12). The observed and calculated wavenumbers in Table 12 and their isotopic shifts in Table 13 are in agreement with the observed vibrational data and hence ClO<sub>4</sub> should exhibit C<sub>2v</sub> symmetry at the minimum of its potential hypersurface.

However, calculated and experimental IR band intensities give in general a much poorer fit than to band positions. The intensity of the ν<sub>8</sub> band is overestimated by 80% and ν<sub>2</sub> by about 250%. The calculated intensities of the bands ν<sub>3</sub>, ν<sub>4</sub>, ν<sub>6</sub>, ν<sub>7</sub> deviate over a greater range but in average they fit well to the experimental intensities. The greatest deviation exhibits the ν<sub>9</sub> band, which has only 10–20% of the calculated intensity.

A particular feature of the ClO<sub>4</sub> IR matrix spectrum is the broadness of the absorption bands with half widths of up to 20 cm<sup>-1</sup>. At higher resolution (0.13 cm<sup>-1</sup>) a fine structure of the band profile appears, as shown for example in Figure 13 d at the ν<sub>6</sub> band. At least six components can be seen and we assume that they are caused by matrix site splittings, which are unusually pronounced in ClO<sub>4</sub> (normally the half width of IR matrix bands are 0.1–1 cm<sup>-1</sup>). The reason for this feature may be the Jahn–Teller effect in ClO<sub>4</sub> resulting in a distortion from T<sub>d</sub> symmetry and a shallow potential minimum.<sup>[133, 134]</sup> In this situation, the ClO<sub>4</sub> radicals “feel” the different matrix cages more than an ordinary molecule and the superposition of all spectra for all different matrix sites (here >6) leads to broad bands. Furthermore, the theoretical calculations point out, that a second stable structure of ClO<sub>4</sub> with C<sub>s</sub> symmetry lies only ≈ 4–5 kJ mol<sup>-1</sup> above the C<sub>2v</sub> minimum.<sup>[133, 134]</sup> The ClO<sub>4</sub> radicals are produced at about 500 K and with regard to the Boltzmann distribution a noticeable amount of ClO<sub>4</sub> with the C<sub>s</sub> structure should be present in the heated spray on

Table 11. Vibrational wavenumbers [cm<sup>-1</sup>]<sup>[a]</sup> of ClO<sub>4</sub> isolated in Ne, Ar and O<sub>2</sub> matrix.

Ne matrix <sup>[b]</sup>		I <sup>[c]</sup>	Ar matrix			Assignment /description of mode <sup>[d]</sup>	O <sub>2</sub> Matrix				Approx. description of mode	
<sup>35</sup> Cl <sup>16</sup> O <sub>4</sub>	<sup>37</sup> Cl <sup>16</sup> O <sub>4</sub> <sup>[b]</sup>		<sup>35</sup> Cl <sup>18</sup> O <sub>4</sub>	<sup>35</sup> Cl <sup>16</sup> O <sub>4</sub>	<sup>37</sup> Cl <sup>16</sup> O <sub>4</sub> <sup>[b]</sup>		C <sub>2v</sub> symmetry		<sup>35</sup> Cl <sup>16</sup> O <sub>4</sub> site A	<sup>35</sup> Cl <sup>16</sup> O <sub>4</sub> site B		<sup>37</sup> Cl <sup>16</sup> O <sub>4</sub> site A
		≈ 2280 br				2ν <sub>8</sub>	A <sub>1</sub>					
≈ 2243	≈ 2220	≈ 1				2ν <sub>1</sub>	A <sub>1</sub>					
	≈ 1811	≈ 1	≈ 1727			ν <sub>8</sub> +ν <sub>6</sub>	A <sub>2</sub>					
	≈ 1760 br	≈ 1	≈ 1690 br		≈ 1750	ν <sub>8</sub> +ν <sub>3</sub>	B <sub>2</sub> <sup>[e]</sup>					
	≈ 1695 br	≈ 3	≈ 1620 br		≈ 1685	ν <sub>1</sub> +ν <sub>3</sub>	A <sub>1</sub>					
	≈ 1480 br	≈ 3	≈ 1410 br		≈ 1475	ν <sub>1</sub> +ν <sub>7</sub>	A <sub>1</sub>	1525.4	1528.8			50 ν (O–O)
	≈ 1457 br	≈ 1	≈ 1390 br		≈ 1450	ν <sub>2</sub> +ν <sub>6</sub>	B <sub>1</sub>					
	≈ 1412 br	≈ 2	≈ 1345 br		≈ 1405	ν <sub>2</sub> +ν <sub>3</sub>	A <sub>1</sub>					
1292.5	1287.7	7	1252		≈ 1285	ν <sub>1</sub> C <sub>s</sub> ClO <sub>4</sub> <sup>[f]</sup>						
1233.6	1221.6	100	1196.4	1227	1217	ν <sub>8</sub> ν <sub>as</sub> (Cl <sup>18</sup> O <sub>2</sub> )	b <sub>2</sub>	1252.1	1244.7	1238	1234	100 ν <sub>as</sub> (ClO <sub>3</sub> )
1207 br		≈ 3				ν <sub>6</sub> +ν <sub>3</sub>	B <sub>1</sub>	1229.4	1196.6	1216.6	1185	
1160.8	1150.4	100	1118.2	1154	1144	ν <sub>1</sub> ν <sub>i</sub> (Cl <sup>18</sup> O <sub>2</sub> ) a <sub>1</sub>						
1132 br		≈ 3				2ν <sub>3</sub>	A <sub>1</sub>	1012.7	1002.3	1010.5	1000	24 ν <sub>s</sub> (ClO <sub>3</sub> )
873.0		4	831.3	≈ 872		ν <sub>2</sub> ν <sub>i</sub> (Cl <sup>16</sup> O <sub>2</sub> )	a <sub>1</sub>					
644.3	642.1	16	608.4	641	639	ν <sub>6</sub> ν <sub>as</sub> (Cl <sup>16</sup> O <sub>2</sub> )	b <sub>1</sub>	687.8br	≈ 681.5	667.0	660.5	58 ν (Cl–O)
573.0	569.9	16	545.5	570	564	ν <sub>3</sub> δ (Cl <sup>18</sup> O <sub>2</sub> )	a <sub>1</sub>					
544.3	540.9	2	518.8			ν <sub>9</sub> ω/ρ(ClO <sub>2</sub> )	b <sub>2</sub>	593.6	589.4		585	15 δ <sub>as</sub> (ClO <sub>3</sub> )
354.0	350.5	3	345.6			ν <sub>7</sub> δ(Cl <sup>16</sup> O <sub>2</sub> )	a <sub>1</sub>	573.5	576.3	569		
337.2	(334.8)	1	321.2			ν <sub>4</sub> ω/ρ (ClO <sub>2</sub> )	b <sub>1</sub>					

[a] Peak position at the main matrix site. [b] Shoulders or peaks at the <sup>35</sup>ClO<sub>4</sub> isotopomer band, localized from the second derivative spectrum. [c] Relative infrared intensities integrated over all matrix sites and isotopomers. [d] Cl<sup>18</sup>O<sub>2</sub> short Cl–O bonds, Cl<sup>16</sup>O<sub>2</sub> long Cl–O bonds. [e] Also possible ν<sub>8</sub>+ν<sub>9</sub>A<sub>1</sub>, ν<sub>1</sub>+ν<sub>6</sub>B<sub>1</sub>. [f] Or 2ν<sub>6</sub> of C<sub>2v</sub> ClO<sub>4</sub> in Fermi resonance with ν<sub>8</sub>, see text.

Table 12. Calculated vibrational wavenumbers [ $\text{cm}^{-1}$ ], intensities ( $\text{km mol}^{-1}$ , in *italic*) bond lengths [pm] and bond angles [ $^\circ$ ] of  $^{35}\text{Cl}^{16}\text{O}_4$  ( $C_{2v}$ ).

Mode	B3LYP/ cc-pVQZ <sup>[a]</sup>		UB3LYP <sup>[b]</sup>		TZ2P <sup>++</sup> B3P86 <sup>[c]</sup>		TZ2P <sup>++</sup> B3LYP <sup>[c]</sup>		cc-pVDZ UHF-CCSD(T) <sup>[d]</sup>		exptl <sup>[e]</sup>	
$\nu_8$ b <sub>2</sub>	1233	<i>(183)</i>	1257	<i>(182)</i>	1229	<i>(190)</i>	1263	<i>(186)</i>	1257	<i>(189)</i>	1234	(100)
$\nu_1$ a <sub>1</sub>	1137	<i>(150)</i>	1160	<i>(150)</i>	1132	<i>(157)</i>	1161	<i>(151)</i>	1161	<i>(157)</i>	1161	(100)
$\nu_2$ a <sub>1</sub>	879	<i>(17)</i>	904	<i>(18)</i>	868	<i>(17)</i>	878	<i>(18)</i>	850	<i>(12)</i>	873	(4)
$\nu_6$ b <sub>1</sub>	623	<i>(25)</i>	648	<i>(22)</i>	615	<i>(20)</i>	612	<i>(18)</i>	609	<i>(12)</i>	644	(16)
$\nu_3$ a <sub>1</sub>	561	<i>(27)</i>	572	<i>(27)</i>	550	<i>(26)</i>	562	<i>(24)</i>	547	<i>(34)</i>	573	(16)
$\nu_9$ b <sub>2</sub>	541	<i>(17)</i>	550	<i>(15)</i>	532	<i>(15)</i>	555	<i>(15)</i>	531	<i>(26)</i>	544	(2)
$\nu_7$ b <sub>1</sub>	400	<i>(3)</i>	413	<i>(2)</i>	397	<i>(2)</i>	366	<i>(6)</i>	450	<i>(4)</i>	354	(3)
$\nu_4$ a <sub>1</sub>	379	<i>(2)</i>	385	<i>(1.7)</i>	374	<i>(2)</i>	383	<i>(2)</i>	391	<i>(2)</i>	337	(1)
$\nu_5$ a <sub>2</sub>	367	<i>(0)</i>	375	<i>(0)</i>	355	<i>(0)</i>	362	<i>(0)</i>	347	<i>(0)</i>	–	
$r$ (Cl–O <sup>s</sup> )	141.4		141.44		142.4		143.3		146.6			
$r$ (Cl–O <sup>l</sup> )	148.8		148.62		150.2		151.6		155.2			
$\angle$ (O <sup>s</sup> ClO <sup>s</sup> )	114.0		113.97		114.2		114.3		114.3			
$\angle$ (O <sup>l</sup> ClO <sup>l</sup> )	94.6		94.75		93.6		93.5		91.4			

[a] Ref. [40]. [b] Ref. [39]. [c] Ref. [133], B3LYP is scaled with 1.07. [d] Ref. [134], scaled with 1.07. [e] This work, relative integrated band intensities in *italics*.

nozzle and should be quenched without relaxation in the matrix. The calculated vibrational wavenumbers for  $\text{ClO}_4$  in  $C_{2v}$  and  $C_s$  symmetry are similar and may overlap each other. However, the most intensive and short-waved band is

predicted to occur in the range of 1300 or 1260  $\text{cm}^{-1}$  for  $\text{Cl}^{16}\text{O}_4$  or  $\text{Cl}^{18}\text{O}_4$ , respectively.<sup>[133]</sup> In Figure 13a such an additional band ( $\approx 5\text{--}7\%$  of  $\nu_8$ ) appears in this region, so that an assignment of this band to  $\text{ClO}_4$  in the  $C_s$  form can be

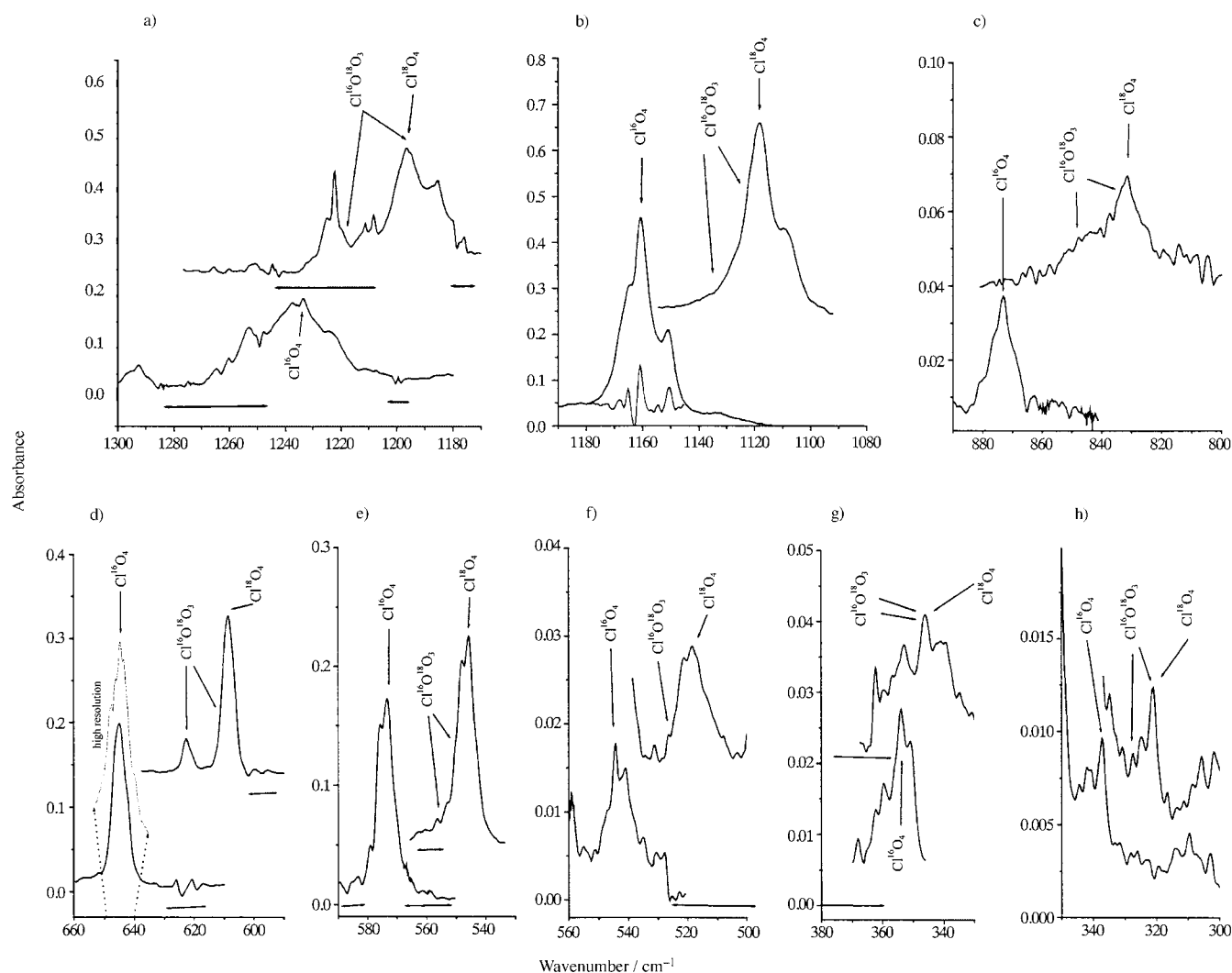


Figure 13. IR band regions of  $\text{Cl}^{16}\text{O}_4$ ,  $\text{Cl}^{18}\text{O}_4$  and  $\text{Cl}^{16}\text{O}^{18}\text{O}_3$  in expanded scale. Arrows indicate calculated band positions (see text). Horizontal arrows indicate by-products subtraction regions

Table 13. Calculated and experimental vibrational isotopic shifts ( $\text{cm}^{-1}$ ) of  $\text{ClO}_4$ .

mode	Exptl wave-number	$\Delta\nu^{35/37}\text{Cl}^{16}\text{O}_4$		$\Delta\nu^{35}\text{Cl}^{16/18}\text{O}_4$			
		calcd <sup>[a]</sup>	exptl <sup>[b]</sup>	calcd <sup>[c]</sup>	calcd <sup>[d]</sup>	calcd <sup>[d]</sup>	exptl <sup>[b]</sup>
$\nu_8$ b <sub>2</sub>	1234	14.8	12	41.3	40	38	37
$\nu_1$ a <sub>1</sub>	1161	11.6	10.2	43.7	41	40	43
$\nu_2$ a <sub>1</sub>	873	1.8	–	48.5	47	44	42
$\nu_6$ b <sub>1</sub>	644	0.1	–	36.4	35	33	36
$\nu_3$ a <sub>1</sub>	573	3.1	3.1	26.6	25	26	28
$\nu_9$ b <sub>2</sub>	544	3.3	3.7	24.6	24	24	26
$\nu_7$ b <sub>1</sub>	354	6.3	3.5	8.7	8	7	8.4
$\nu_4$ a <sub>1</sub>	337	0.4	( $\approx 2$ )	21.3	21	20	16

[a] This work, B3LYP/6-311G (d,p). [b] Shoulders or peaks at the  $^{35}\text{ClO}_4$  isotopomer band, localized from the second derivative spectrum. [c] Ref. [39], UB3LYP. [d] Ref. [133], TZ2P++B3P86 and TZ2P++B3LYP.

taken into account. Hence, an additional reason for the broadness of the bands is seen in an easy distortion of the  $C_{2v}$  structure in direction towards of the  $C_s$  structure.

For evaluation of the  $^{35/37}\text{ClO}_4$  isotopic shifts (see Table 13) the second derivative spectra are used as shown for the  $\nu_1$  band in Figure 13 b (this band is undisturbed by precursor and by-product bands). Three groups of two bands separated by  $10.2\text{ cm}^{-1}$  ( $\Delta\nu^{35/37}\text{ClO}_4$ ) are clearly seen.

The IR spectrum of  $^{18}\text{O}$  isotopically labelled  $\text{ClO}_4$ , which contains about 15%  $\text{Cl}^{16}\text{O}^{18}\text{O}_3$ , possesses another interesting feature. For an estimate of the band positions of the  $\text{Cl}^{16}\text{O}^{18}\text{O}_3$  isotopomer, we calculated the wavenumbers of the  $\text{Cl}^{16}\text{O}_4$ ,  $\text{Cl}^{18}\text{O}_4$  ( $C_{2v}$ ), and  $\text{Cl}^{16}\text{O}^{18}\text{O}_3$  ( $C_s$ , with  $^{16}\text{O}$  located at the short or long bond) fundamentals at the B3LYP/6-311G(p,d) level of DFT theory.<sup>[216]</sup> The results for the scaled wavenumbers of  $\text{Cl}^{16}\text{O}^{18}\text{O}_3$  are gathered in Tables 14 and the band positions are

Table 14. Scaled calculated vibrational wavenumbers [ $\text{cm}^{-1}$ ] of the  $\text{Cl}^{16}\text{O}^{18}\text{O}_3$  in  $C_s$  symmetry based on the calculated and experimental vibrational wavenumbers for  $\text{Cl}^{16}\text{O}_4$  and  $\text{Cl}^{18}\text{O}_4$  in  $C_{2v}$  symmetry.

mode	$\text{Cl}^{16}\text{O}_4$ , $C_{2v}$		$\text{Cl}^{18}\text{O}_4$ , $C_{2v}$		factor <sup>[b]</sup>	$\text{Cl}^{16}\text{O}^{18}\text{O}_3$ , $C_s$ short <sup>[c]</sup>		$\text{Cl}^{16}\text{O}^{18}\text{O}_3$ , $C_s$ long <sup>[c]</sup>	
	exptl	calcd <sup>[a]</sup>	exptl	calcd <sup>[a]</sup>		scaled <sup>[d]</sup>	scaled <sup>[d]</sup>		
$\nu_8$ b <sub>2</sub>	1234	1149	1196	1112	1.075	a'	1219	a''	1195
$\nu_1$ a <sub>1</sub>	1161	1055	1118	1017	1.100	a'	1134	a'	1122
$\nu_2$ a <sub>1</sub>	873	798	831	755	1.097	a'	835	a'	846
$\nu_6$ b <sub>1</sub>	644	543	608	512	1.187	a''	611	a'	622
$\nu_3$ a <sub>1</sub>	573	521	546	497	1.099	a'	557	a'	550
$\nu_9$ b <sub>2</sub>	544	504	519	481	1.079	a'	526	a''	526
$\nu_7$ b <sub>1</sub>	354	313	346	306	1.131	a''	347	a'	348
$\nu_4$ a <sub>1</sub>	337	358	321	338	0.945	a'	321	a'	327

[a] This work B3LYP/6-311G (d,p). [b] Scaling factor: exptl/calcd averaged value. [c]  $^{16}\text{O}$  placed on the short or long bond. [d] This work B3LYP scaled with the factor from [b].

indicated by arrows in Figure 13. All band profiles of  $\text{Cl}^{16}\text{O}_4$  and  $\text{Cl}^{18}\text{O}_4$  should be the same, but they are distorted for  $\text{Cl}^{18}\text{O}_4$  in all cases (see Figure 13). With the overlap of  $\text{Cl}^{16}\text{O}^{18}\text{O}_3$  bands one can rationalize the changes in the  $\text{Cl}^{18}\text{O}_4$  band profiles. Only for  $\nu_6$ , one component of  $\text{Cl}^{16}\text{O}^{18}\text{O}_3$  is clearly separated from  $\text{Cl}^{18}\text{O}_4$ . The appearance of this component is a further evidence for the two different ClO bonds in  $\text{ClO}_4$ .

In contrast to the other chlorine oxide radicals,  $\text{ClO}_4$  interacts strongly with oxygen. The typical broad feature of the IR spectrum of matrix isolated  $\text{ClO}_4$  vanishes completely, when oxygen is used as a matrix material. At  $1525\text{ cm}^{-1}$  a strong OO stretching band is seen which is comparable to the analogous band of FOO, isolated in a oxygen matrix ( $1496\text{ cm}^{-1}$ ).<sup>[217]</sup> Four additional groups with sharp IR bands appear near  $1250$ ,  $1010$ ,  $690$  and  $590\text{ cm}^{-1}$ , which are all split into two components (site A and B) as in the OO stretching band. There is an interesting site selectivity by photolysis. After irradiation of the  $\text{O}_2$  matrix with light of wavelength  $>395\text{ cm}^{-1}$  the intensity of site B bands decrease and at shorter photolysis wavelengths both sites disappear and ClOO is observed as photoproduct. All vibrational data of the new species are collected in Table 11. The band positions and relative intensities as well as the  $^{35/37}\text{Cl}$  isotopic shift of the latter four bands are very similar to the spectrum of  $\text{FCIO}_3$  ( $\nu_1$ ,  $\Delta\nu^{35/37}\text{Cl}$ , int.:  $1316$ ,  $19$ , vs;  $1062$ ,  $3$ , m;  $712$ ,  $10$ , s;  $588$ ,  $3\text{ cm}^{-1}$ , m). Hence, the new radical  $\text{O}_3\text{ClO-O}_2$  has been formed, which is comparable to the recently described isoelectronic species  $\text{FSO}_2\text{O-O}_2$ .<sup>[218]</sup> Both new peroxy radicals will be investigated in detail in a future matrix study.

**The UV/Vis spectrum of  $\text{ClO}_4$ :** Our experimental setup allows the measurement of IR and UV/Vis spectra from the same matrix, which guaranties that the same absorber correlates in both spectral regions. The corresponding UV/Vis spectra of the five best low pressure flash thermolysis experiments of  $\text{Cl}_2\text{O}_6$  are used to “extract” the UV spectrum of  $\text{ClO}_4$  from the raw data. For this purpose reference spectra of  $\text{Cl}_2\text{O}_6$ ,  $\text{ClO}_3$  and  $\text{ClO}_2$  are used to subtract their content in the raw spectra. Finally the raw spectra are averaged and smoothed in the region  $370$  to  $420\text{ nm}$ , because in this region incompenations by subtraction of  $\text{ClO}_2$  bands appeared. Below  $370\text{ nm}$  no reliable absorption profile can be evaluated, due to strongly absorbing by-products. The resulting spectrum is shown in Figure 14.

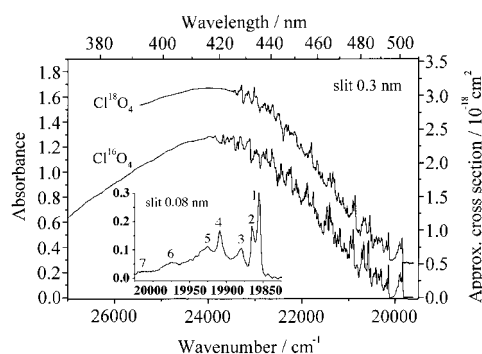


Figure 14. UV/Vis spectrum of  $\text{ClO}_4$  isolated in a neon matrix. The spectrum of  $\text{Cl}^{18}\text{O}_4$  is shifted up by 0.3 absorbance units. Bands of  $\text{ClO}_2$  are subtracted in the region  $\lambda > 420\text{--}430\text{ nm}$ . For  $\lambda < 420\text{--}430\text{ nm}$  the spectrum is smoothed (see text). In the insert a high resolution spectrum of the onset is shown.

In Figure 15 an expanded scale spectrum including the spectrum of the  $^{18}\text{O}$  isotope enriched species are shown. For the subtraction of the by-products  $\text{Cl}^{16/18}\text{O}_2$  and  $\text{Cl}^{16/18}\text{O}_3$  separate thermolysis experiments with  $\text{Cl}_2^{16/18}\text{O}_6$  at higher



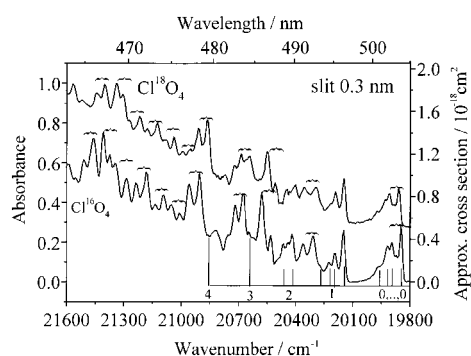


Figure 15. UV/Vis spectrum of  $\text{ClO}_4$  isolated in a neon matrix in expanded scale, showing the overlying vibrational bands. Brackets indicate peaks shifted by about  $50 \text{ cm}^{-1}$  (see text, Table 15). The assignment of some peaks to a progression is shown by the vertical lines.

temperatures are performed, so that the amount of by-products [see Eqs. (7.2), (41)–(43)] is enhanced. The resulting  $\text{Cl}^{18}\text{O}_4$  spectrum is somewhat disturbed by  $\text{Cl}^{16}\text{O}^{18}\text{O}_3$  isotopomers (see above). The absorption cross section scale in Figures 14 and 15 is estimated from the  $\text{ClO}_2$  absorbance at  $420.58 \text{ nm}$  ( $4,0,0$  transition) in the raw spectra and its known cross section<sup>[23]</sup> assuming a band width similar to that of gaseous  $\text{ClO}_2$  at room temperature, and assuming a 1 : 1 mixture of  $\text{ClO}_2$  :  $\text{ClO}_4$  according to Equation (7.2).

Above  $510 \text{ nm}$  no absorption is observed, which is in accordance with the photoelectron spectrum of the  $\text{ClO}_4^-$  anion<sup>[25]</sup> where no electron detachments between 1 and  $2.5 \text{ eV}$  ( $20000 \text{ cm}^{-1}$ ) above the  $X^2B_1$  ground state are detectable. The UV/Vis spectrum of  $\text{ClO}_4$  shows a strongly structured band with a maximum at about  $23500 \text{ cm}^{-1}$ , a half width of about  $5000 \text{ cm}^{-1}$  and a band origin at  $19847 \text{ cm}^{-1}$  for  $\text{Cl}^{16}\text{O}_4$ . The reproducible fine structure on the long wavelength wing is most unusual on three accounts: i) the  $0-0$  transition line consists of seven components as displayed in the insert of in Figure 14; ii) the  $0-0$  transition of  $\text{Cl}^{16}\text{O}_4$  is about  $12 \text{ cm}^{-1}$  lower in energy, than that of  $\text{Cl}^{18}\text{O}_4$  (uncertainty of the difference is  $\pm 5 \text{ cm}^{-1}$ ) and iii) the fine structure is very complex.

Splittings of bands due to electronic transition are not unexpected in the matrix. They can amount to  $30-70 \text{ cm}^{-1}$  as measured for  $\text{FCO}_2$ <sup>[159]</sup> in a Ne matrix or for  $\text{ClO}_2$ <sup>[190]</sup> in a Ar matrix. However, the separations between the bands in the insert in Figure 14 are  $10, 24, 55, 70, 117, \text{ and } 155 \text{ cm}^{-1}$  ( $0.08 \text{ nm}$  slit). As discussed in the preceding section, a portion of  $\text{ClO}_4$  in the  $C_s$  form can be recognized; the two satellites ( $117, 155 \text{ cm}^{-1}$ ) can be caused by some  $\text{ClO}_4$  distorted to the less stable  $C_s$  structure in two matrix cages. Furtheron, the separation between the first and the fourth band ( $\approx 50 \text{ cm}^{-1}$ ,  $0.3 \text{ nm}$  slit) in the  $0-0$  transition seem to be the main feature in the fine structure, as it appears at other band groups, also indicated by brackets in Figure 15. Unfortunately, the contour of the  $0-0$  transition is nearly faded out in the rest of the spectrum. Only the first progression can be analysed in some detail, supported by its contour ( $1,0$ ; and  $2,0$ ) and further bands from the second derivative spectrum. This transition can be analysed up to the quantum number 4 (see Figure 15, Table 15) and a quadratic fit results in Equations (47) and (48).

Table 15. Positions of some progression bands of Ne matrix isolated  $\text{ClO}_4$  in the  $A^2B \leftarrow X^2A$  transition.

Site	$\text{Cl}^{16}\text{O}_4$ wavenumber/ $\text{cm}^{-1}$	$\text{Cl}^{18}\text{O}_4$ wavenumber/ $\text{cm}^{-1}$	Assignment
1,2	19847	19859	0,...,0
4	19895	19906	0,...,0
5	19918	19930	0,...,0
6	19972	19972	0,...,0
1,2	20147	20147	1,0
4	20193	20190	1,0
5	20227	20214	1,0
6	20270	20257	1,0
1,2	20415	20400	2,0
4	20469	20450	2,0
5	20494	20475	2,0
1,2	20640sh	20601sh	3,0
1,2	20840sh	20807sh	4,0
1,2	20307	20288	0,1 ?
4	20362	20356	0,1 ?
5	20387sh		0,1 ?
1,2	20672	20640	0,2 ?
4,5	20717	20684	0,2 ?
1,2	20905	20859	0,3 ?
4,5	20964	20908	0,3 ?
1,2	21090	21037	0,4 ?
4,5	21137	21077	0,4 ?
1,2	21182	21124	0,5 ?
4,5	21238	21177	0,5 ?
1,2	21286	21218	?
4,5	21340	21265	?
1,2	21456	21402	?
4,5	21512	21445	?
1,2	20532	20507	inverted intensity
4,5	20576	20547	inverted intensity
1,2	21011	20955	inverted intensity
4,5	21050	20990	inverted intensity
1,2	21374	21299	inverted intensity
4,5	21409	21340	inverted intensity

$$\text{Cl}^{16}\text{O}_4: \nu = 19847 (\pm 1) + 320 (\pm 1) \cdot N - 18 (\pm 0.3) \cdot N^2 \quad (47)$$

$$\text{Cl}^{18}\text{O}_4: \nu = 19859 (\pm 1) + 304 (\pm 1) \cdot N - 16 (\pm 0.2) \cdot N^2 \quad (48)$$

Theoretical calculations examine a  $^2B_1$  ground state and an  $^2A_2$  excited state<sup>[134]</sup> so that all  $a_1$  and  $b_2$  vibrations are allowed as progression bands.<sup>[212]</sup> Taking the relative isotopic shift of  $\text{Cl}^{16/18}\text{O}_4$  into account the  $\nu_2$  ( $a_1$ ) vibration belongs to this progression, which is lowered by 64% in the excited state. This vibrational motion is described as the symmetric stretching of the longer Cl–O bonds. Furthermore, the intensity is strongly decreased (analyzed by subtracting the totally smoothed spectrum from the original spectrum), which point to a continuum at quantum numbers higher than five.

Assignment of the other transitions is difficult to perform and must be tentative due to the possibility of combined progression. The data for progression bands up to  $21500 \text{ cm}^{-1}$  are gathered in Table 15 and taken from the second derivative spectrum. In comparison to the UV/Vis spectrum of  $\text{ClO}_4$  the spectrum of isoelectronic  $\text{SO}_4^-$ <sup>[219, 220]</sup> is structureless and red-shifted by about  $25 \text{ nm}$  with a cross section of  $6.3 \times 10^{-18} \text{ cm}^2$  at  $440-450 \text{ nm}$ , which is about three times greater than our estimated cross section of  $\text{ClO}_4$ . Furthermore a high energy band of lower intensity is observed at about  $320 \text{ nm}$  in the UV/Vis spectrum of  $\text{SO}_4^-$ . In this region our spectrum is strongly perturbed by  $\text{ClO}_2$  and/or  $(\text{OClO})_2$ .

**The Photochemistry of ClO<sub>4</sub>:** Light of wavelengths > 495 nm is sufficient to deplete Ne matrix isolated ClO<sub>4</sub>. The only detectable IR active product is OClO and therefore the photolysis should proceed as follows [Eq. (49)].



It is assumed that this photodissociation is a one-step process in the  $A^2A_2 \leftarrow X^2B_1$  transition where the symmetric stretching of the longer ClO bonds is active, strongly decreased in energy and follow the proper coordinate.

ClO<sub>4</sub> isolated in an oxygen matrix seems to be more photostable. Light of wavelengths > 395 nm is needed for the photolysis of the more photolabile matrix site and light of wavelengths < 395 nm is suitable for photolysing both matrix sites. In both cases secondary photoisomerization of OClO to ClOO occurs.

## Conclusion

The spectroscopic, bonding, structural and chemical properties of the radicals ClO<sub>x</sub> ( $x = 1-4$ ) are very unusual in many respects. Especially ClO<sub>4</sub> stands out among these species. Due to a strong Jahn–Teller effect two of the ClO bonds are stretched leading to a  $C_{2v}$  equilibrium configuration. The  $C_s$  distorted geometry is found to be only a few kJ mol<sup>-1</sup> higher in energy. The results of several recent ab initio calculations<sup>[38–40, 133, 134]</sup> are in accordance with the observed IR spectrum of ClO<sub>4</sub>. A reason for the unusual broad IR bands is seen in the shallow potential which makes the ClO<sub>4</sub> radical floppy. Also the photochemistry of ClO<sub>4</sub> differs from that of the other radicals. While ClO<sub>x</sub> ( $x = 1-3$ ) photodissociate, primarily into ClO<sub>x-1</sub>+O the ClO<sub>4</sub> radical yields ClO<sub>2</sub>+O<sub>2</sub>. Irregularities of the properties of the chlorine oxide radicals with increasing oxidation numbers are seen in: i) the energy onset for photodissociation follows the sequence ClO<sub>3</sub> < ClO<sub>4</sub> < ClO<sub>2</sub> < ClO and consequently ClO<sub>3</sub> is the least and ClO the most photostable species; ii) dimerization leads for ClO to different unstable Cl<sub>2</sub>O<sub>2</sub> isomers, for OClO to weakly bonded dimers, for ClO<sub>3</sub> to stable asymmetric O<sub>2</sub>ClOClO<sub>3</sub> and for ClO<sub>4</sub> to decomposition products (presumably to Cl<sub>2</sub>O<sub>6</sub> + O<sub>2</sub>); iii) ClO bond lengths and bond energies follow the sequence ClO < ClO<sub>3</sub> < ClO<sub>2</sub> and ClO<sub>3</sub> < ClO<sub>2</sub> < ClO, respectively. The only nearly even trend appears in the electron affinity (eV): ClO (2.5),<sup>[174]</sup> ClO<sub>2</sub> (2.15), ClO<sub>3</sub> (4.25), ClO<sub>4</sub> (5.25).<sup>[25]</sup> The very high electron affinity of ClO<sub>4</sub> may be one reason for its strong interaction with molecular oxygen under formation of the peroxy radical O<sub>3</sub>ClO-O<sub>2</sub>. However, the electron affinity of ClO<sub>4</sub> is apparently not sufficient to form a stable O<sub>2</sub>+ClO<sub>4</sub><sup>-</sup> salt, because the ionization energy of dioxygen amounts 12.07 eV.<sup>[174]</sup>

In contrast, in the series of the anions ClO<sub>x</sub><sup>-</sup> ( $x = 1-4$ ), which contain just one electron more than the respective radicals, their bonding properties are changing smoothly with increasing oxidation number. Their bond lengths (pm) and force constants (10<sup>2</sup> N m<sup>-1</sup>) gradually changes in the series ClO<sup>-</sup> (-/3.3), ClO<sub>2</sub><sup>-</sup> (157/4.3), ClO<sub>3</sub><sup>-</sup> (149/5.87), ClO<sub>4</sub><sup>-</sup> (145/7.45)<sup>[199, 221]</sup> and their bond angles can be rationalized by the VSEPR rules.

## Experimental Section

**Caution:** Chlorine oxides and ozone are potentially explosive, especially in the presence of oxidizable materials. It is important to take proper safety precautions when these compounds are handled and reactions are carried out. Reactions involving either one of them should be carried out only with millimolar quantities.

**General procedures and reagents:** Volatile materials were manipulated in a glass vacuum line equipped with two capacitance pressure gauges (221 AHS-1000 and 221 AHS-10 MKS Baratron, Burlington, MA), three U-traps, and valves with PTFE stems (Young, London, UK). The vacuum line was connected to an IR cell (optical path length 200 mm, Si windows 0.5 mm thick) contained in the sample compartment of a FTIR instrument. This allows one to observe the purification processes and to follow the course of reactions. The compounds ClO<sub>3</sub>,<sup>[222]</sup> HOClO<sub>3</sub>,<sup>[223]</sup> ClOClO<sub>3</sub>,<sup>[224]</sup> FOClO<sub>3</sub>,<sup>[225]</sup> ClO<sub>2</sub>ClO<sub>4</sub>,<sup>[47, 86]</sup> and ClONO<sub>2</sub>,<sup>[226]</sup> were prepared according to literature procedures and stored in flame-sealed glass ampoules under liquid nitrogen in a storage Dewar vessel. By using an ampoule key<sup>[227]</sup> the ampoules were opened at the vacuum line, an appropriate amount was taken out for the experiments and then they were flame-sealed again. <sup>18</sup>O-labelled ClO<sub>2</sub> was prepared by treating ClF<sub>3</sub> with D<sub>2</sub><sup>18</sup>O.<sup>[228]</sup> Isotopically enriched <sup>18</sup>O<sub>3</sub> was synthesized from <sup>18</sup>O<sub>2</sub> in a small home-made ozonizer.<sup>[229, 230]</sup> Excess of <sup>18</sup>O<sub>2</sub> was recovered by cryopumping oxygen into a vessel filled with molecular sieve (5 Å) held at -196 °C. To prepare <sup>18</sup>O-enriched CsClO<sub>4</sub>, CsCl was dissolved in D<sub>2</sub><sup>18</sup>O and oxidized by an excess of XeF<sub>2</sub> first at room temperature and then for 2 h at 80 °C. According to the literature the synthesis of the isotopically labelled Cl<sub>2</sub><sup>18</sup>O<sub>6</sub> is accomplished by the reaction of Cl<sup>18</sup>O<sub>2</sub> and <sup>18</sup>O<sub>3</sub>. Due to incomplete labelled precursors (ClO<sub>2</sub>: 93% Cl<sup>18</sup>O<sub>2</sub>, 6% Cl<sup>16</sup>O<sup>18</sup>O, 1% Cl<sup>16</sup>O<sub>2</sub>; O<sub>3</sub>: 94% <sup>18</sup>O<sub>3</sub>, 6% other isotopomers) up to 20% of <sup>16/18</sup>O<sub>2</sub>Cl<sup>16</sup>OCl<sup>18</sup>O<sub>3</sub> and other isotopomers were produced.

**Chemicals:** The following chemicals were obtained from commercial sources and used without purification: H<sub>2</sub>SO<sub>4</sub> (96%), HClO<sub>4</sub> (73%), KClO<sub>3</sub>, KClO<sub>4</sub>, CsClO<sub>4</sub> (all p.a.: Merck, Darmstadt, Germany) D<sub>2</sub><sup>18</sup>O (98.97% <sup>18</sup>O, Johnson Matthey, Alfa Products, Karlsruhe, Germany), F<sub>2</sub> (commercial grade, Solvay, Hannover, Germany), <sup>18</sup>O<sub>2</sub> (99.5%, Chemtrade, Düsseldorf, Germany), O<sub>2</sub> (99.999%, Linde, München, Germany), Ar (99.999%, Messer Griesheim, Krefeld, Germany) and Ne (> 99.999%, Messer Griesheim, Krefeld, Germany).

**Preparation of the matrices:** Small amounts of the samples (ca. 0.1 mmol) were transferred in vacuo into a small U-trap kept in liquid nitrogen. This U-trap was mounted in front of the matrix support (a metal mirror) and allowed to reach a temperature of -130 °C or -120 °C for ClO<sub>2</sub>, -105 °C for HOClO<sub>3</sub> and -38 °C for Cl<sub>2</sub>O<sub>6</sub> in cold baths. In the temperature range -130 to -110 °C and -110 to -30 °C Dewar vessels filled with isopentane or ethanol were used, respectively. A gas stream (≈ 3 mmol h<sup>-1</sup>) of argon, oxygen or neon was directed over the cold sample in the U-trap, and the resulting gas mixtures passed the heated quartz nozzle (Ø 4 mm with an end orifice of 1 mm) within milliseconds and were quenched on the matrix support at 12 or 5 K, respectively. Because the vapor pressures of the samples in the cold U-trap are of the magnitude of 10<sup>-3</sup> mbar and the pressures of the inert gas streams during the deposition are about 1 mbar in the U-trap, the resulting sample to gas ratios can be estimated to be in the range 1:1000. In case of ClO<sub>2</sub> the higher temperatures enables greater sample to gas ratios and formation of an increased fraction of (ClO<sub>2</sub>)<sub>2</sub> molecules during the quenching process. For each sample two different amounts of matrix material (≈ 1 and 3 mmol, respectively) were deposited through the heated nozzle at temperatures of 200 for ClO<sub>2</sub>, 350–380 for ClOClO<sub>3</sub>, 350 for FOClO<sub>3</sub>, 240 for Cl<sub>2</sub>O<sub>6</sub> 400–450 for HOClO<sub>3</sub> and 420–450 °C for ClONO<sub>2</sub>. Heating the ClO<sub>2</sub>/inert gas stream was necessary for dissociation of evaporated (ClO<sub>2</sub>)<sub>2</sub> molecules. Mixtures of NO<sub>2</sub>: Ne (1:1000) were prepared in a stainless steel high-vacuum line and transferred via a stainless steel capillary to the heated nozzle, and quenched as a matrix at 5 K.

**IR spectroscopy:** Gas-phase IR spectra were recorded with a resolution of 2 cm<sup>-1</sup> in the range of 4000–400 cm<sup>-1</sup>, using a FTIR instrument (Nicolet, Impact 400D, Madison, WI) which was directly coupled with the vacuum line. Matrix IR spectra were recorded on a Bruker IFS 66v/S FT spectrometer (Bruker, Karlsruhe, Germany) in the reflectance mode, using a transfer optic. A DTGS or MCT detector together with a KBr/Ge

beam splitter operated in the regions of 5000–400 or 7000–600  $\text{cm}^{-1}$ , respectively. In these regions 64 scans were co-added for each spectrum using apodized resolutions of 1.0, 0.25 or 0.13  $\text{cm}^{-1}$ . A far IR-DTGS detector together with a Ge coated 6  $\mu\text{m}$  Mylar beam splitter was used in the region of 650–80  $\text{cm}^{-1}$ . In this region 64 scans were co-added for each spectrum using an apodized resolution of 1.0  $\text{cm}^{-1}$ .

**UV/Vis spectroscopy:** UV/Vis spectra were recorded in the region 200–750 nm with a Perkin–Elmer Lambda 900 instrument (Perkin–Elmer, Norwalk, CT) using a slit of 0.3 nm, data point separation of 0.15 nm and an integration time of 0.48 s for each datapoint in the range 200–490 (510 for  $\text{ClO}_4$ ) nm. At higher wavelength a datapoint separation of 0.4 nm was used. For matrix measurements two 2 m long quartz single fibers with a special condenser optic (Hellma, Jena Germany) were used. Mercury and Neon pencil lamps were used for wavelength calibration (Oriel Instruments, Stratford, CT).

**Photolysis experiments:** These were performed on the matrices in the visible and UV regions by using a 250 W tungsten halogen lamp (Osram, München, Germany) or a high-pressure mercury lamp (TQ 150, Heraeus, Hanau, Germany), respectively, in combination with a water cooled quartz condenser optic and cut off or interference filters (Schott, Mainz (Germany)). Details of the matrix apparatus have been described elsewhere.<sup>[159]</sup>

For analysis of the IR and UV/Vis matrix spectra of the pyrolysis products, a series of reference spectra of  $\text{ClO}_2$ ,  $\text{Cl}_2\text{O}_4$ ,  $\text{Cl}_2\text{O}_6$ ,  $\text{FClO}_4$ ,  $\text{HClO}_4$ ,  $\text{ClO}$  and  $\text{NO}_2$  have been measured. For  $\text{FO}_2$  and  $\text{F}_2\text{O}_2$  literature data were used.<sup>[1]</sup>

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